

Future of Axion Searches

CPAD Instrumentation Frontier Meeting

Oct 5th, 2015

Gianpaolo Carosi



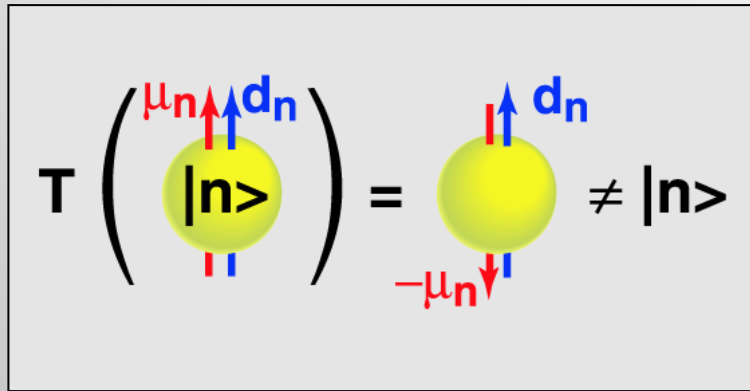
LLNL-PRES-677763

This work was performed under the auspices of the U.S. Department of Energy by Lawrence Livermore National Laboratory under Contract DE-AC52-07NA27344. Lawrence Livermore National Security, LLC



The axion and the Strong-CP problem

QCD CP violation should, e.g., give a large neutron electric dipole moment ($\cancel{T} + CPT = C\cancel{P}$); none is unobserved.
(9 orders-of-magnitude discrepancy)



Why doesn't the neutron have an electric dipole moment?

$$d_e < 3 \cdot 10^{-26} \text{ e-cm}$$

Baker et al. 2006

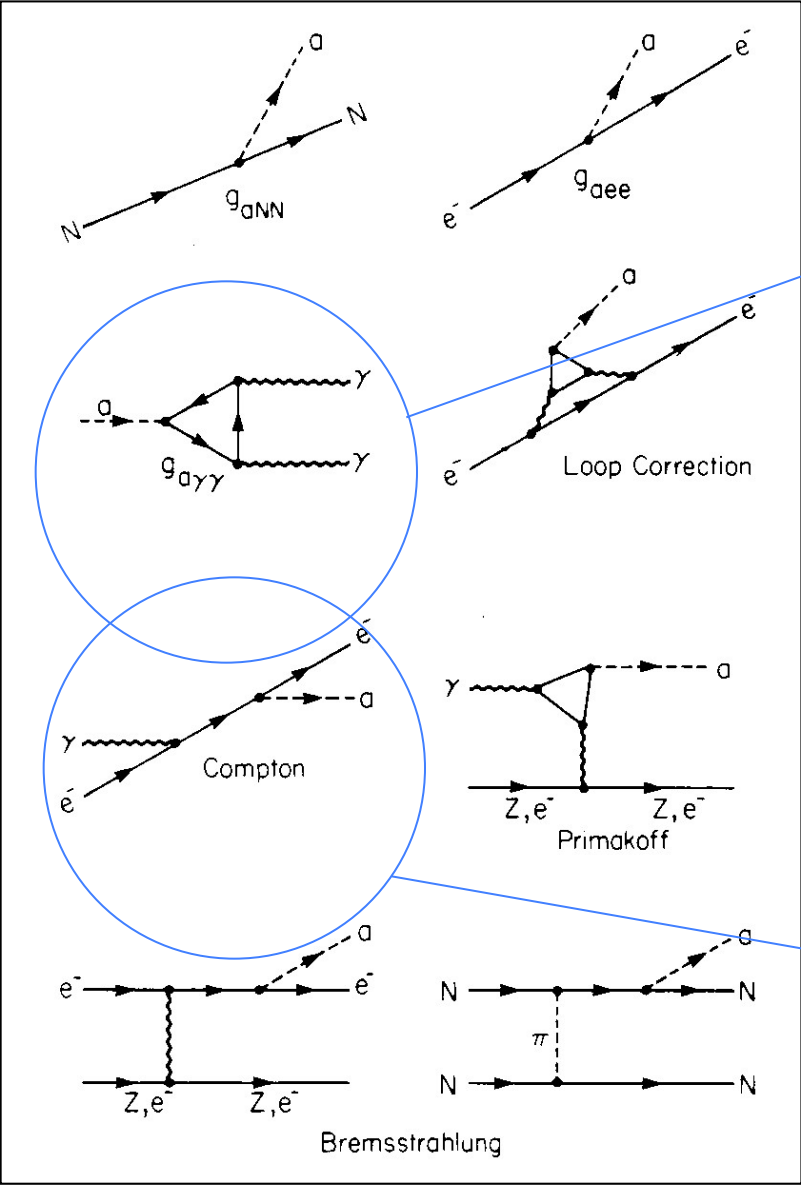
This leads to the “Strong CP Problem”: Where did QCD CP violation go?

1977: Peccei and Quinn: Posit a hidden broken $U(1)$ symmetry \Rightarrow

- 1) A new Goldstone boson (the axion);
- 2) Remnant axion VEV nulls QCD CP violation.

Axion is a pseudoscalar neutron boson (similar to a very light π^0)

Axion Couplings



Coupling to two photons
 Small model uncertainty
 Exploited in certain terrestrial searches
 Easily calculable

Rate depends on “unification group”
 (that is, the particles in the loops),
 ratio of u/d quark masses,
 and mostly f_{PQ}

$$g_{a\gamma\gamma} \sim \frac{\alpha}{f_{PQ}} \left(\frac{E}{N} - 1.95 \right)$$

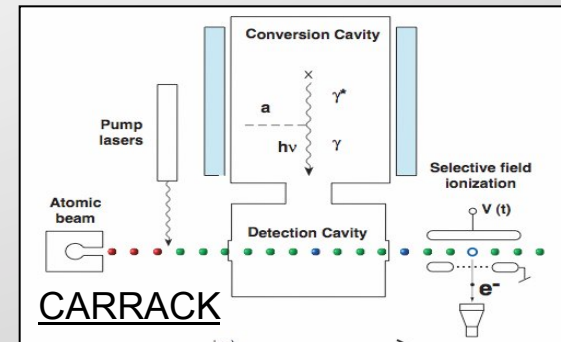
A process with large model uncertainty
 Can occur, e.g., in the Sun
 Contains unknown $U(1)_{PQ}$ charge of electron

Kolb & Turner

Variety of experiments*...

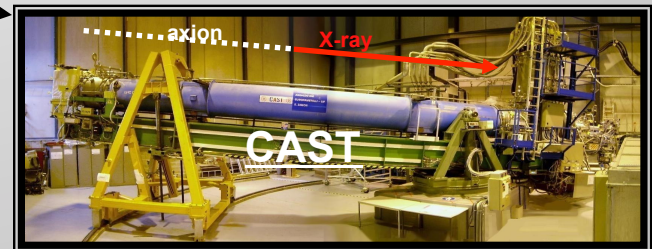
■ Microwave Cavities (dark matter source)

- Low noise amplifiers (**ADMX**) and Rubidium Atoms (**CARRACK**)
 - Look for dark matter axions (low mass) converting to photons in B-Field
 - **Relies on a dense source of primordial axions**



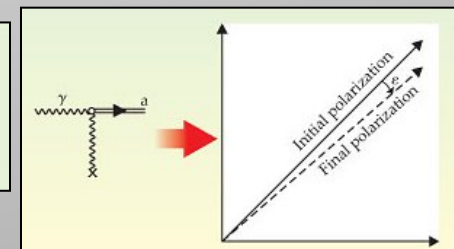
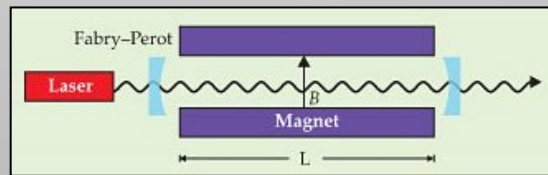
■ Solar Observatories (solar source)

- X-Ray (**CAST**) and Germanium detectors
 - Look for axions generated from the sun
 - Higher coupling required than for DM axions.

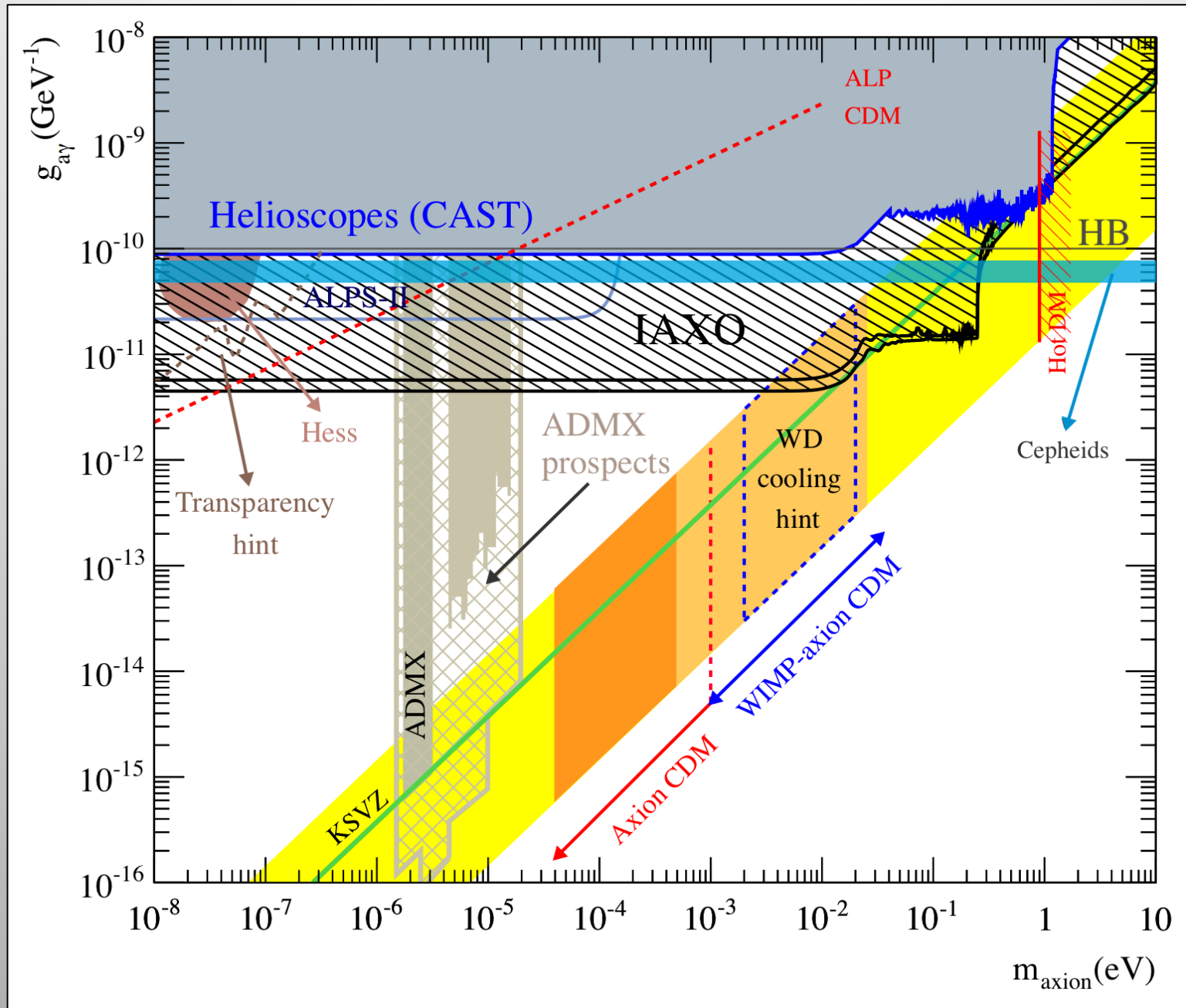


■ Lab experiments (laser source)

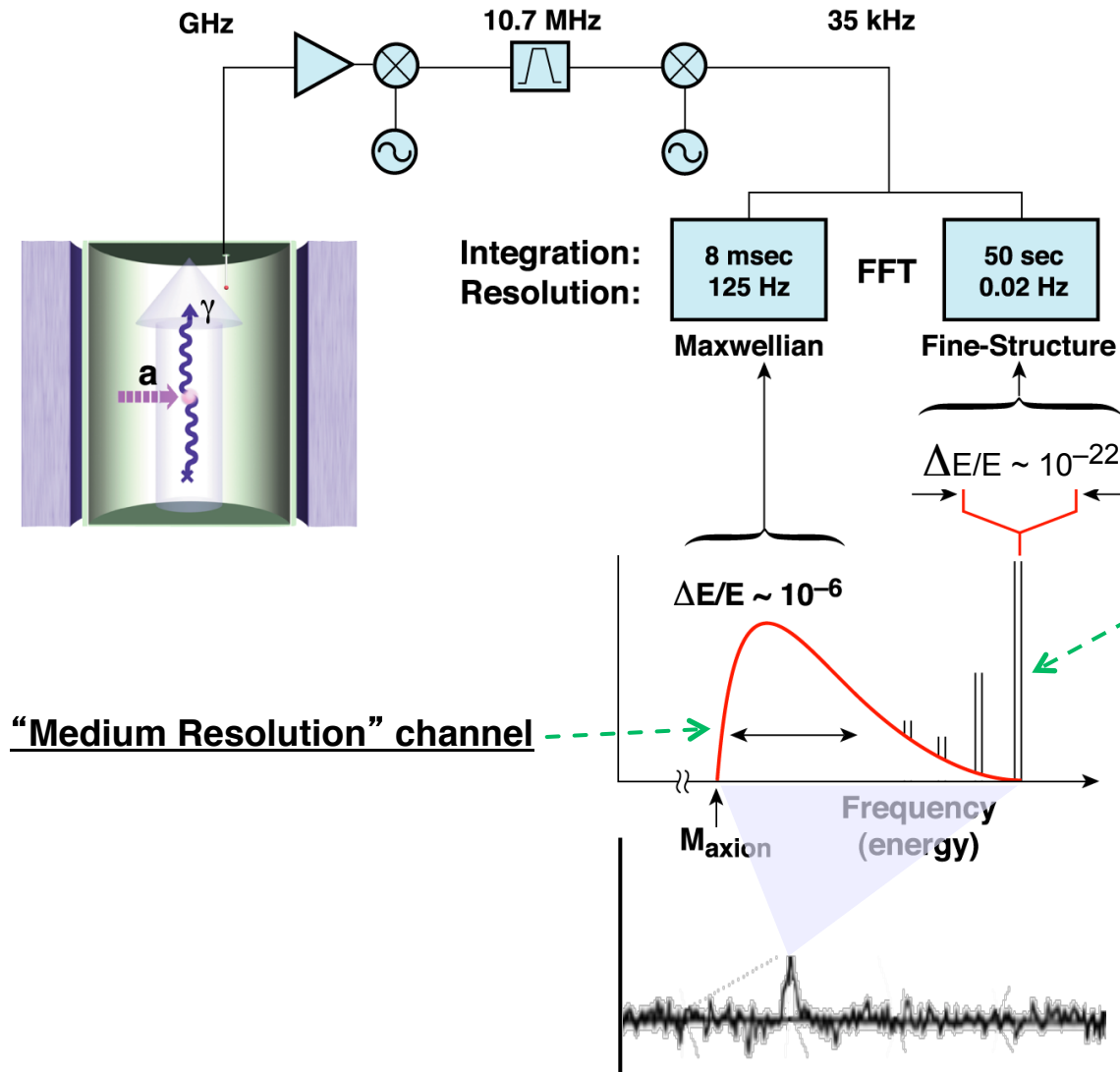
- Photon regeneration and polarization changes (**PVLAS**)
 - Look for production of axions from light passing through B-field
 - Higher coupling required.
 - Ultralight axions (nano-eV) (NMR / LC Circuit)



* See Jeremy Mardon & Surjeet Rajendran talks next



The Axion Dark Matter eXperiment (original concept from P. Sikivie)



Local Milky Way density:

$$\rho_{halo} \sim 450 \text{ MeV/cm}^3$$

Thus for $m_a \sim 10 \mu\text{eV}$:

$$\rho_{halo} \sim 10^{14} \text{ cm}^{-3}$$

“High Resolution” channel

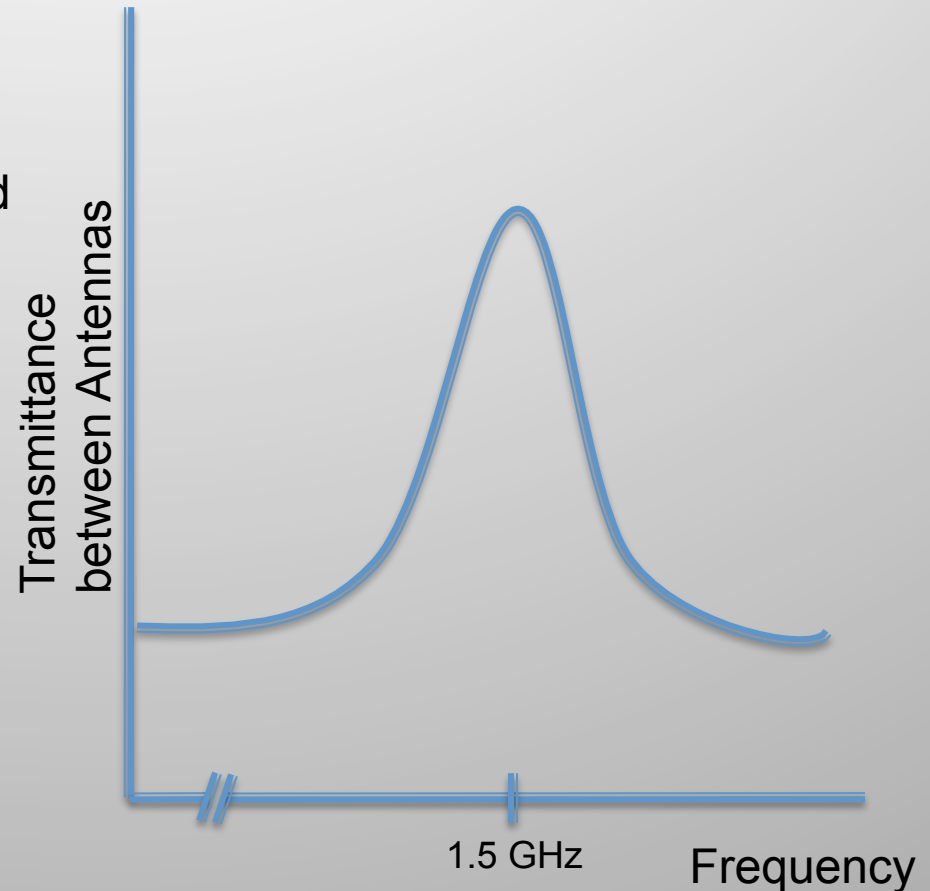
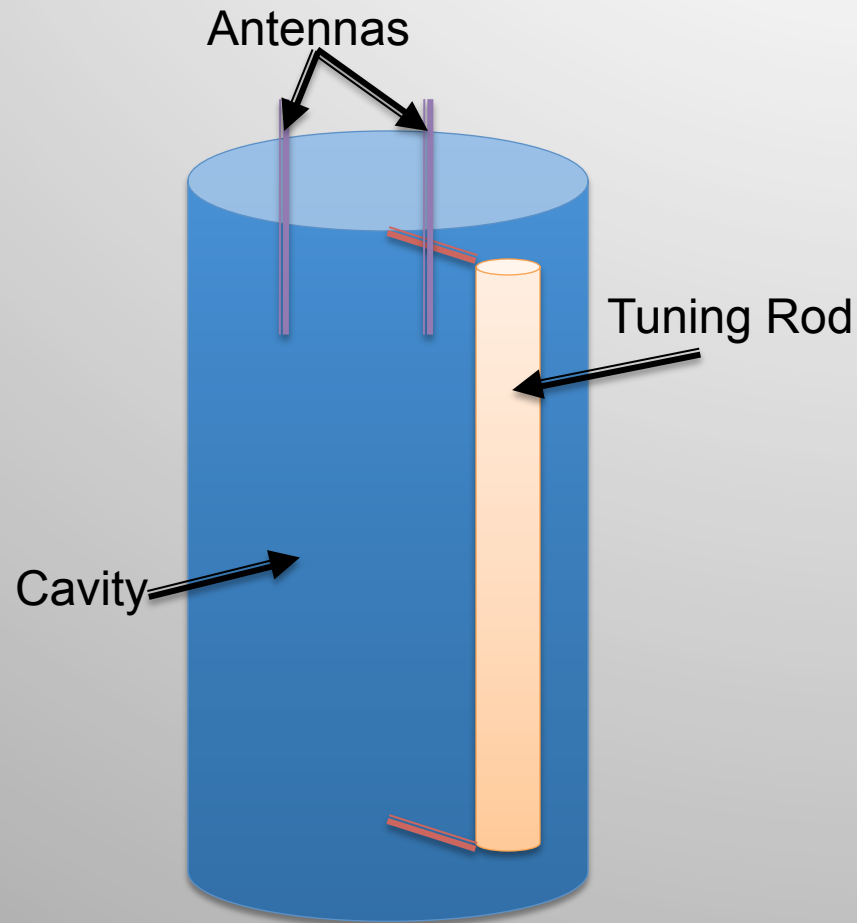
$$\beta_{\text{virial}} \sim 10^{-3} :$$

$$\lambda_{\text{De Broglie}} \sim 100 \text{ m}$$

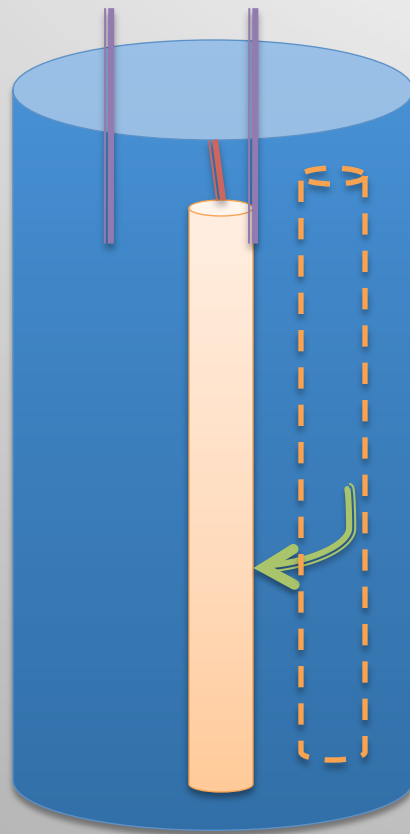
$$\Delta \beta_{\text{flow}} \sim 10^{-11} :$$

$$\lambda_{\text{Coherence}} \sim 1000 \text{ km}$$

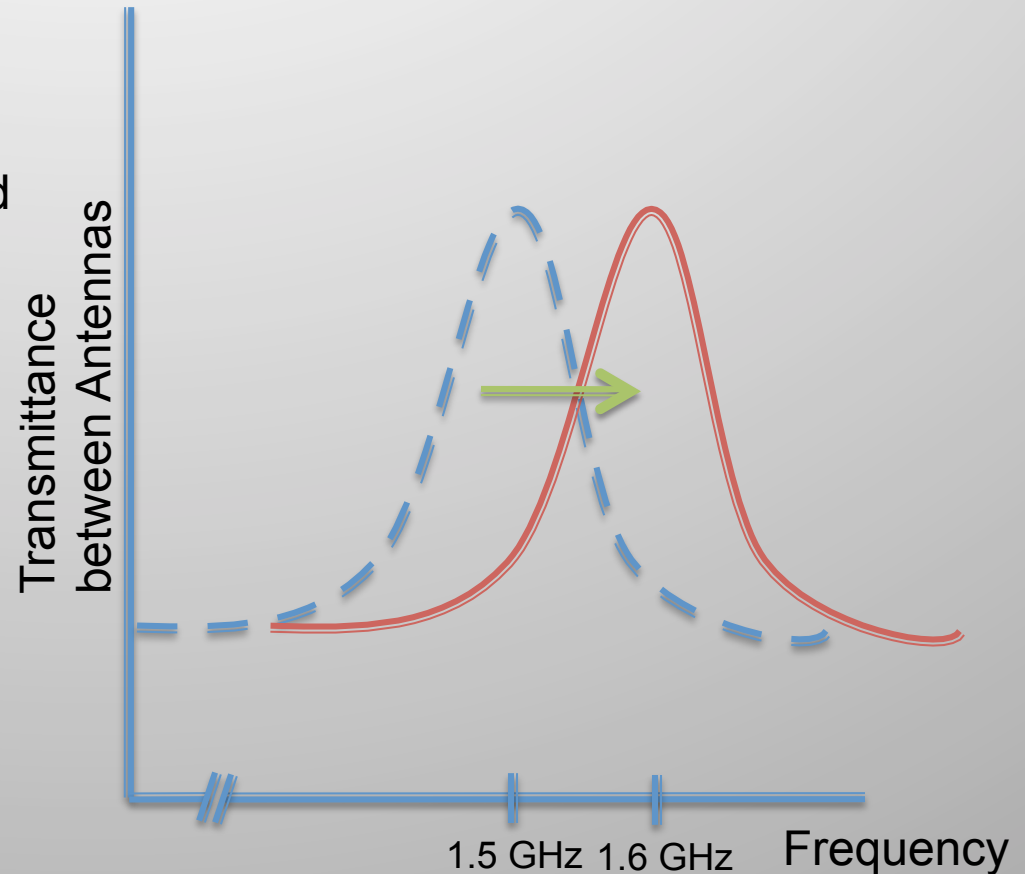
Microwave Cavity needs tunable resonance



Microwave Cavity needs tunable resonance



Tuning Rod



Signal strength

- Power generated in the cavity from axion signal:

$$P_a \approx 3.38^{-23} \text{ W} \left(\frac{V}{100 \text{ l}} \right) \left(\frac{B_0}{8 \text{ Tesla}} \right)^2 \left(\frac{C_{nml}}{0.5} \right) \left(\frac{g_\gamma}{0.36} \right)^2 \cdot \left(\frac{\rho_a}{0.3 \text{ GeV/cc}} \right) \left(\frac{m_a}{1 \text{ GHz}} \right) \left(\frac{\min(Q_L, Q_a)}{10^5} \right)$$

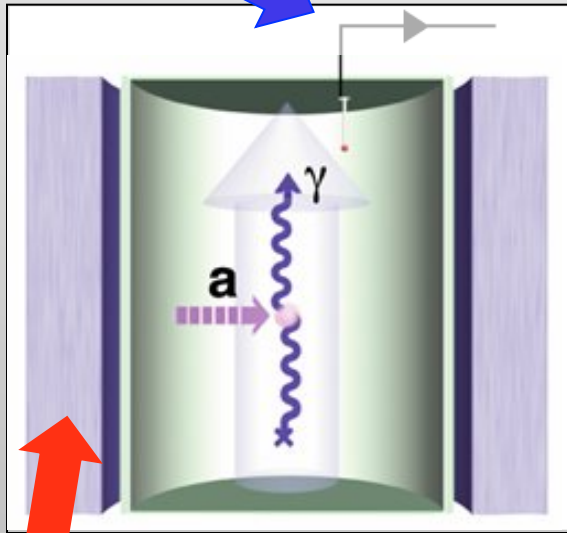
- ρ_a is the local halo density,
- m_a the axion mass (where 1 GHz \sim 4.1 ueV),
- C_{nml} is form factor (overlap of resonant mode E-Field with B-field)
- $Q_L \sim 10^5$ (loaded cavity quality factor); $Q_a \sim 10^6$ (axion line width)
- $g_Y \sim 0.97$ (KSVZ); $g_V \sim 0.36$ (DFSZ) are coupling strengths

The radiometer eqn.* dictates the strategy

$$\frac{s}{n} = \frac{P_{sig}}{kT_S} \cdot \sqrt{\frac{t}{\Delta\nu}}$$

But integration time limited to ~ 100 sec

* Dicke, 1946



System noise temp. now

$$T_S = T + T_N \sim 1.5 + 1.5 \text{ K}$$

But $T_{Quant} \sim 30 \text{ mK}$

This is where we invested to get to Gen 2

$$P_{sig} \sim (B^2 V Q_{cav})(g^2 m_a \rho_a) \\ \sim 10^{-23} \text{ Watts for ADMX}$$

But magnet size, strength $B^2 V \sim \$$

Scan Rate from Dicke Radiometer equation

- Rate determined from SNR

$$\frac{df}{dt} \approx 750 \text{ MHz/year} \left(\frac{g_\gamma}{0.36} \right)^4 \left(\frac{5}{SNR} \right)^2 \left(\frac{f}{1 \text{ GHz}} \right)^2 \left(\frac{B_0}{8 T} \right)^4 \left(\frac{V}{100 l} \right)^2 \left(\frac{Q_L}{10^5} \right) \left(\frac{C_{010}}{0.5} \right)^2 \left(\frac{0.2 \text{ K}}{T_{sys}} \right)^2$$

- SNR is the Signal-to-Noise for detection (usually set to 5),
- f is the frequency being searched (where $1 \text{ GHz} \sim m_a = 4.1 \text{ ueV}$)
- T_{sys} is the total system temperature ($T_{sys} = T_{cavity} + T_{amps}$)

The Importance of Low Noise Temperature

- Original system noise temperature: $T_S = T + T_N = 3.2 \text{ K}$
Cavity temperature: $T = 1.5 \text{ K}$ (pumped ^4He)
Amplifier noise temperature: $T_N = 1.7 \text{ K}$ (HEMT)
- Time* to scan the frequency range from $f_1 = 0.5$ to $f_2 = 1 \text{ GHz}$:

$$\tau(f_1, f_2) = 4 \times 10^{17} (3.2\text{K}/1 \text{ K})^2 (1/f_1 - 1/f_2) \text{ sec} \approx \mathbf{130 \text{ years}}$$

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- Next generation (ADMX Gen 2):
Cavity temperature: $T = 100 \text{ mK}$ (^3He dilution unit)
Amplifier noise temperature: $T_N = 50 \text{ mK}$ (MSA)

*Dine-Fischler-Srednicki-Zhitnitsky (DFSZ) theory

The Importance of Low Noise Temperature

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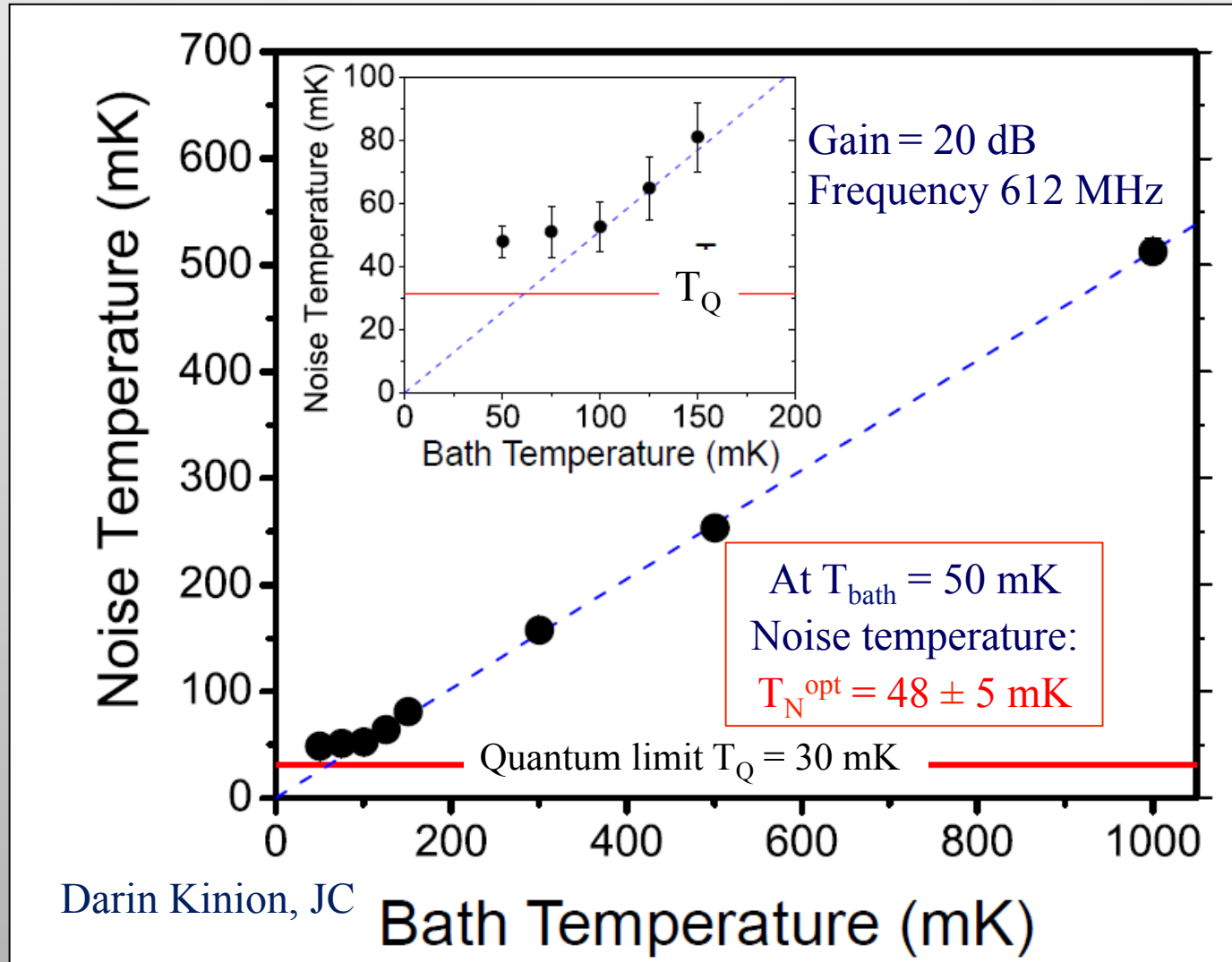
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- Next generation:
Cavity temperature: $T = 100 \text{ mK}$ (^3He dilution unit)
Amplifier noise temperature: $T_N = 50 \text{ mK}$ (MSA)
- Time* to scan the frequency range from $f_1 = 0.5$ to $f_2 = 1 \text{ GHz}$:

$$\tau(f_1, f_2) = 4 \times 10^{17} (0.15\text{K}/1 \text{ K})^2 (1/f_1 - 1/f_2) \text{ sec} \approx \mathbf{104 \text{ days}}$$

*Dine-Fischler-Srednicki-Zhitnitsky (DFSZ) theory

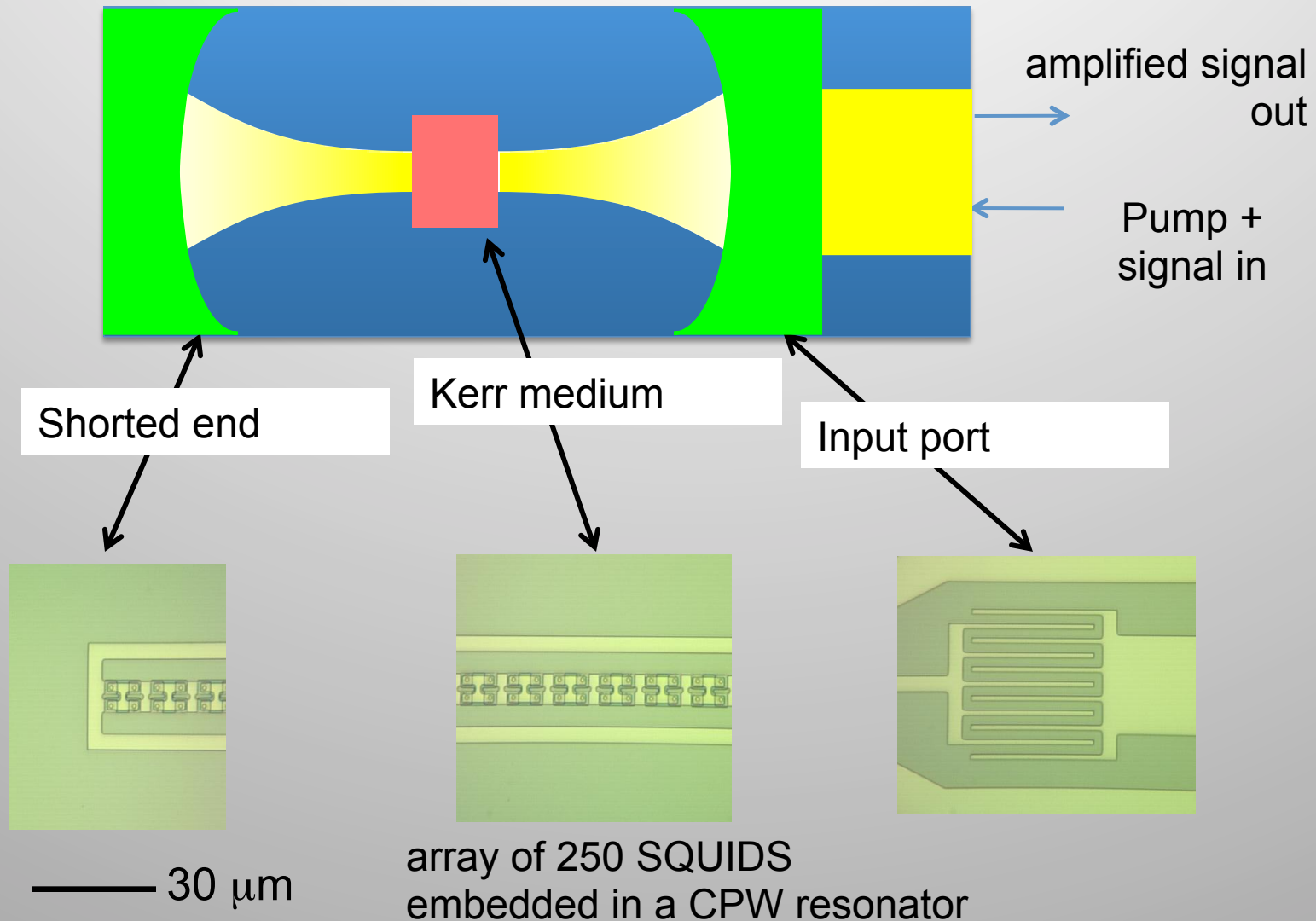
Enabling Technology: Microstrip SQUID Amps



Noise temperatures of 48 ± 5 mK have been demonstrated at 612 MHz, within 1.7 times the quantum limit

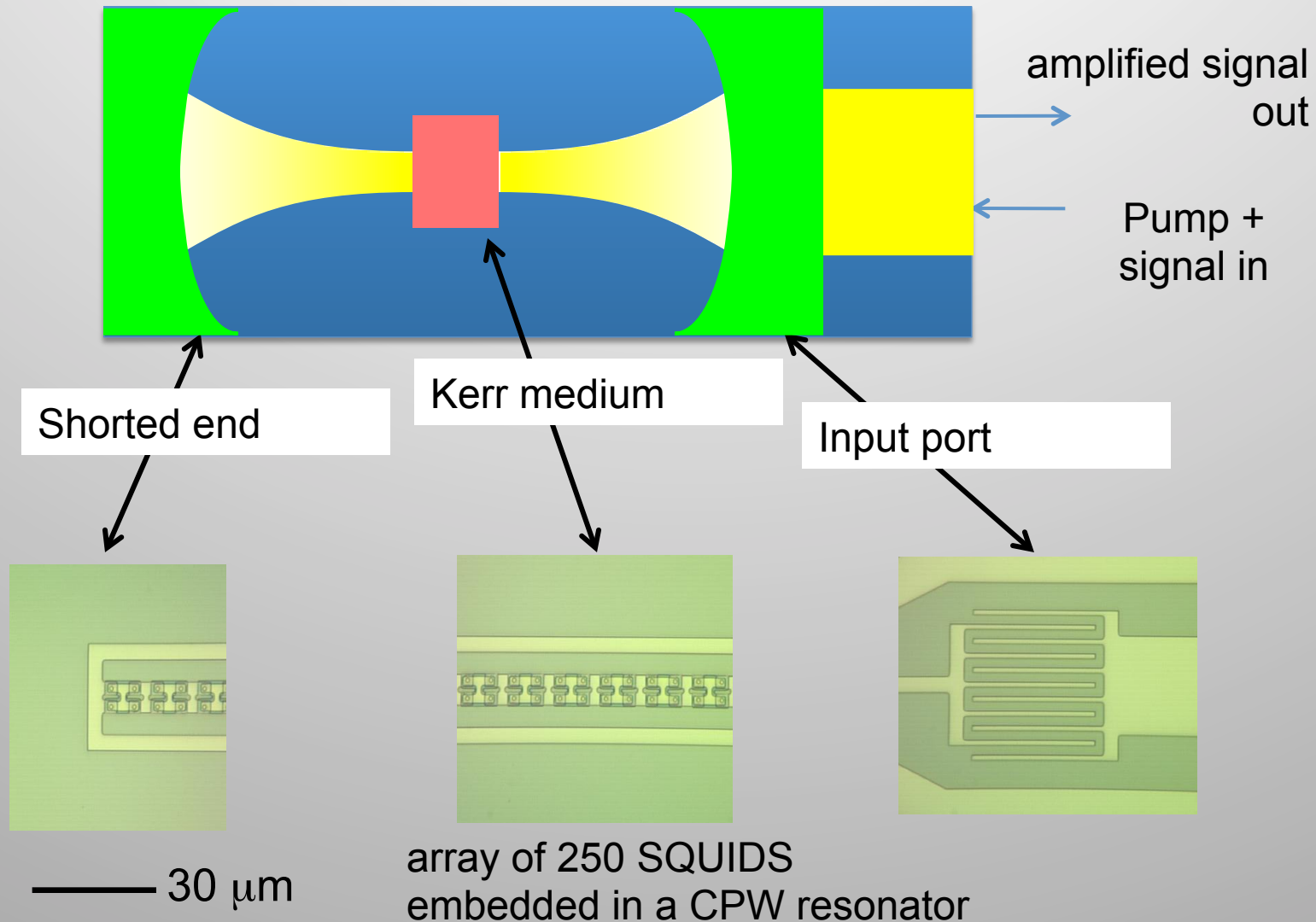
Josephson Parametric Amplifiers (JPA):

SQUIDs embedded in a cavity form a parametric amplifier

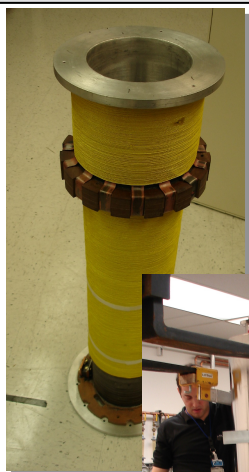


Josephson Parametric Amplifiers (JPA):

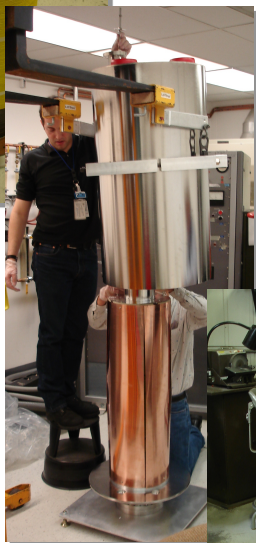
SQUIDs embedded in a cavity form a parametric amplifier
(see Saptarshi Chaudhuri's quantum amp talk tomorrow)



ADMX Experimental Apparatus

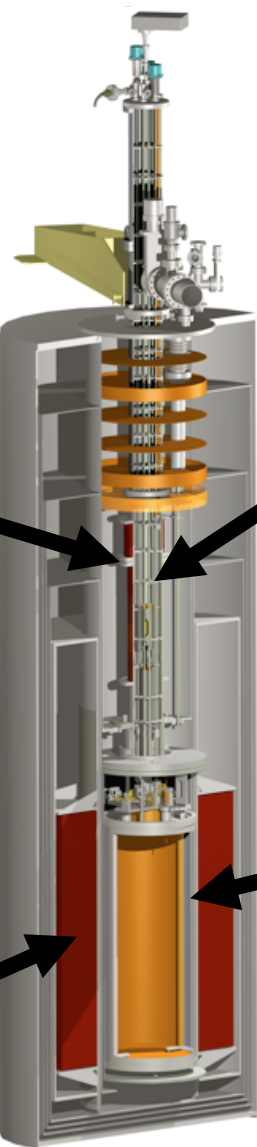


Field compensation magnet for SQUIDs

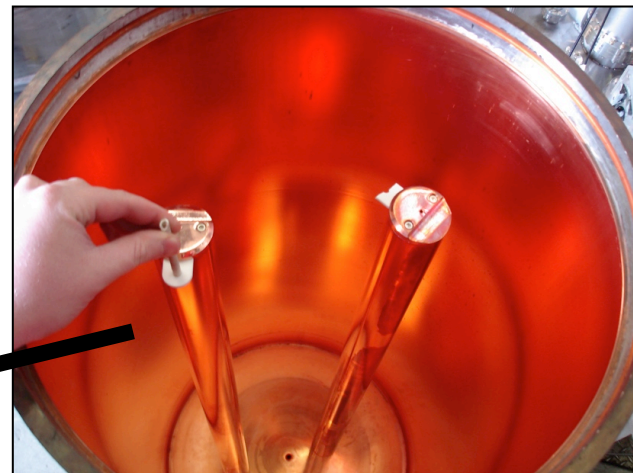
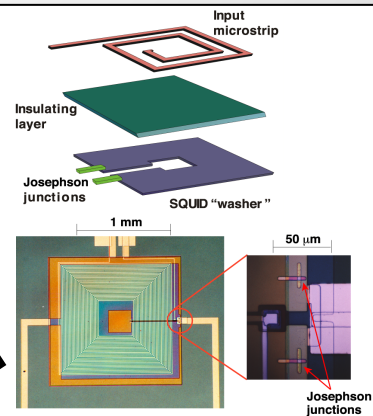


8 Tesla Magnet

11'

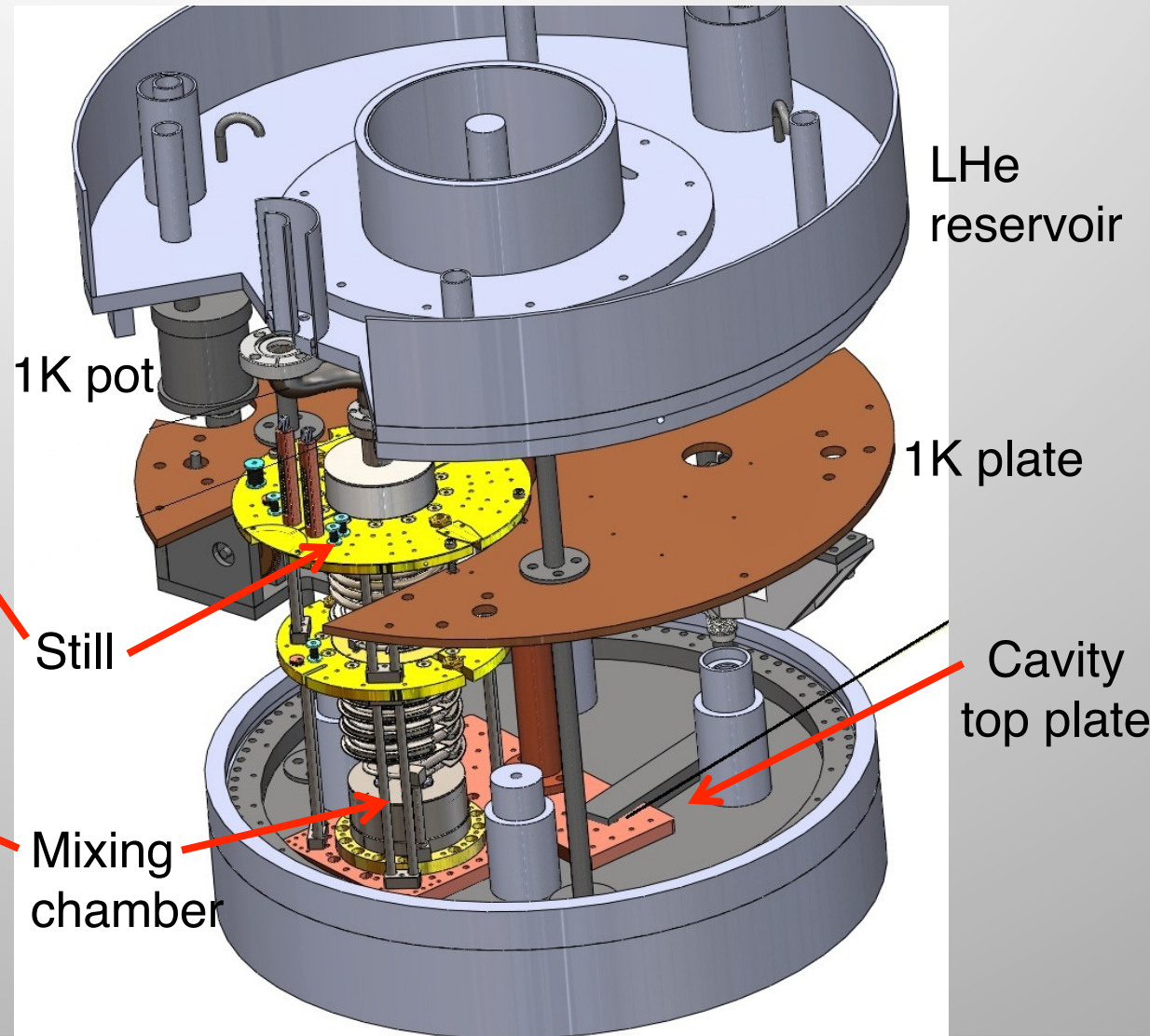
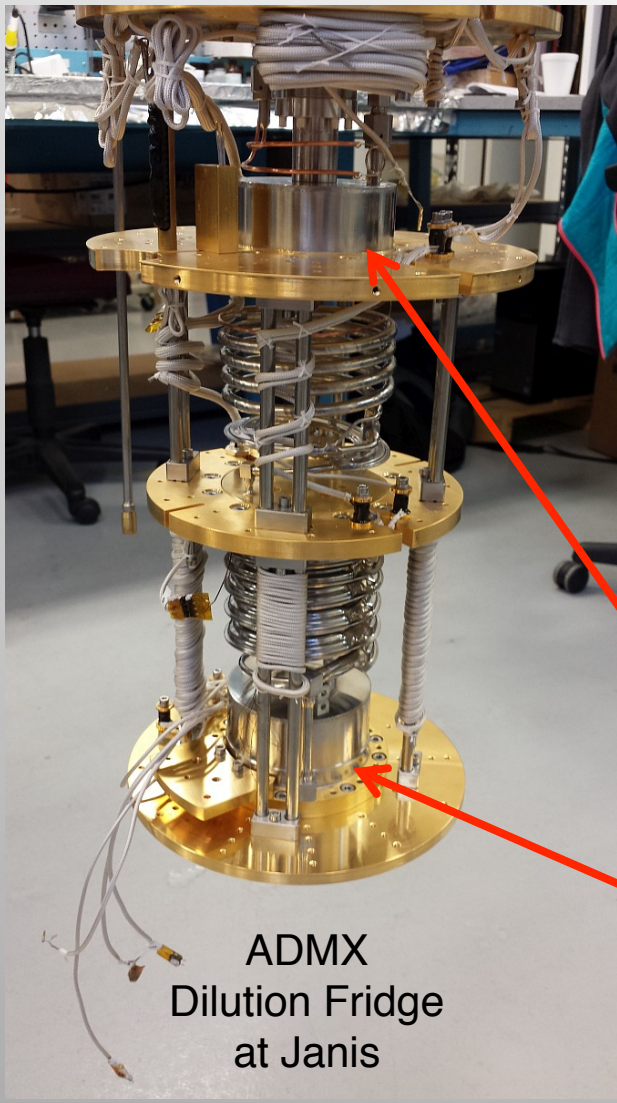


SQUID amplifier

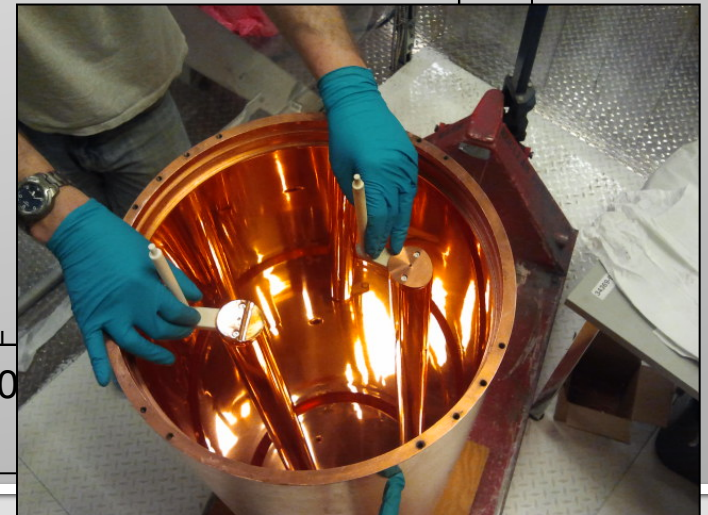
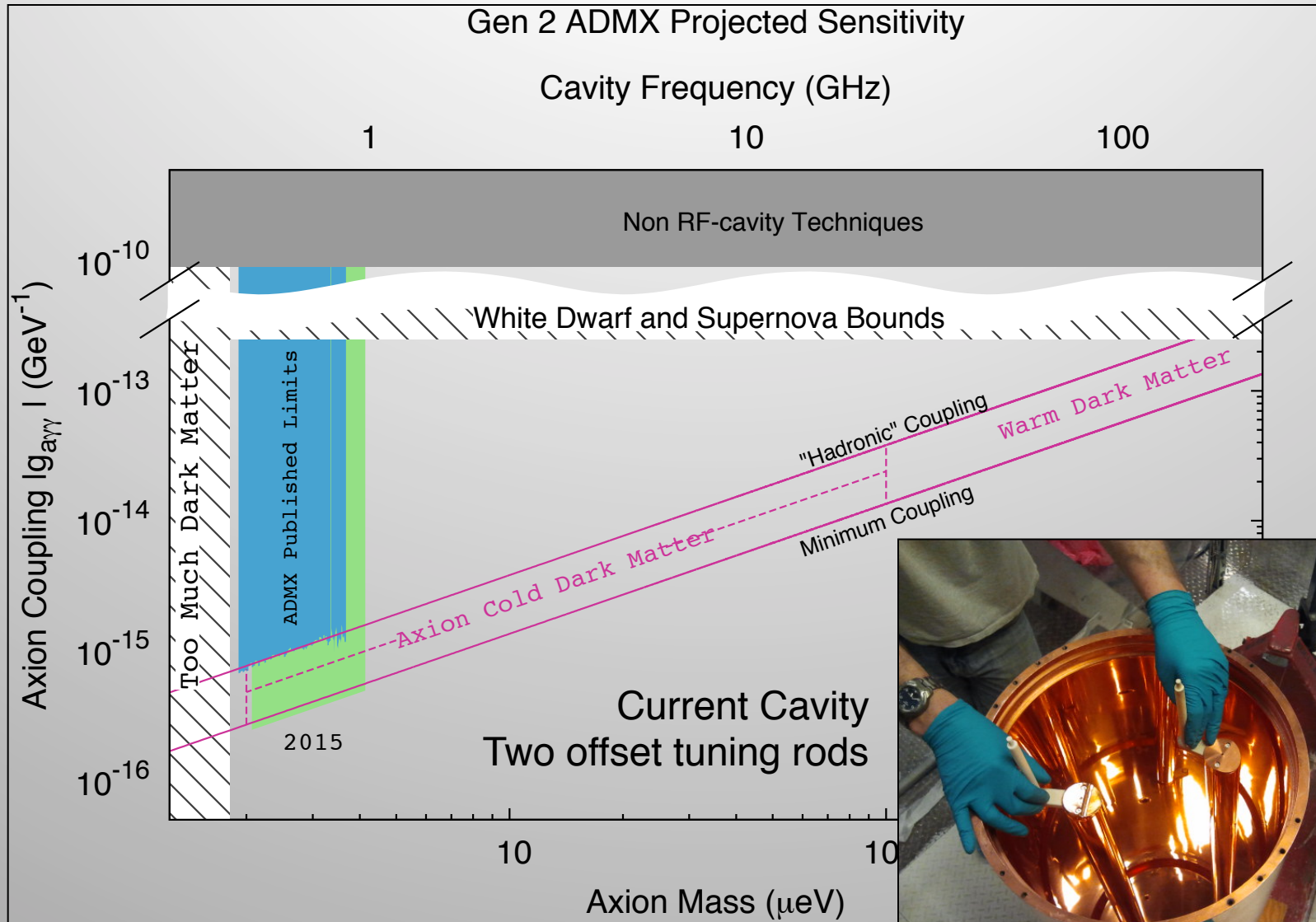


*140 liter microwave cavity
(500 MHz - 1 GHz)*

Dil. Fridge (800 μ W at 100 mK) being commissioned now



ADMX Gen 2 Science Prospects: Year 1 (0.5 – 1 GHz)

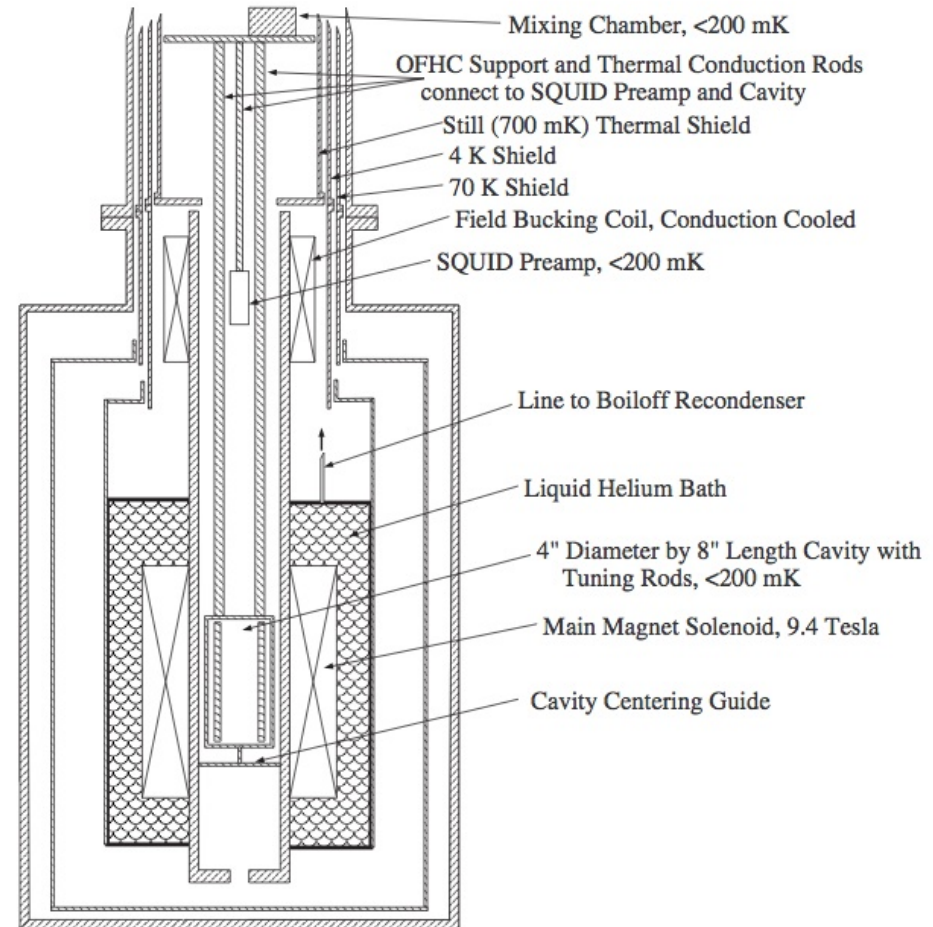


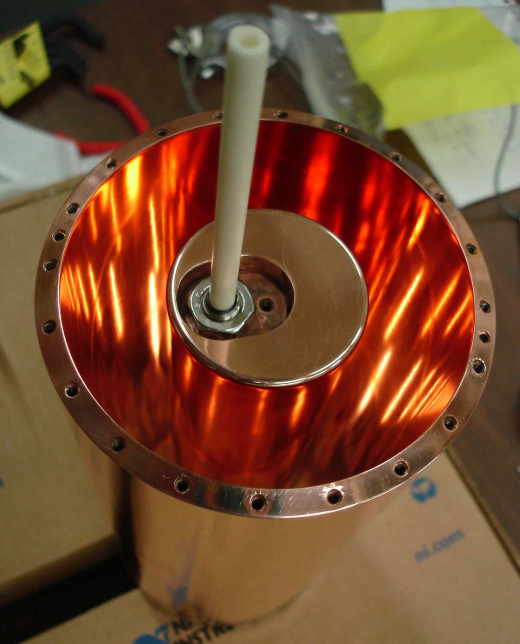
ADMX – HF: High Frequency

Second ADMX site: Yale University – NSF Sponsored

PI: Prof. Steve Lamoreaux

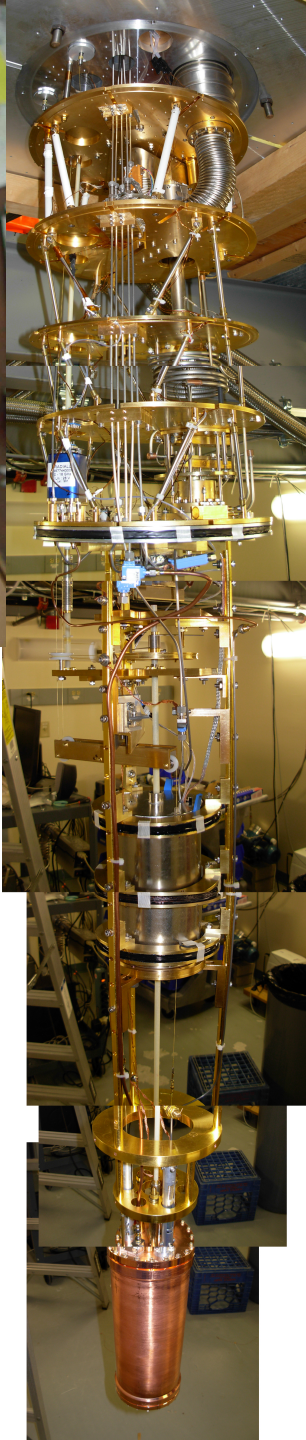
- **New Superconducting Magnet**
5" diameter, 20" long, 9.4 T
- **Existing Dilution fridge.**



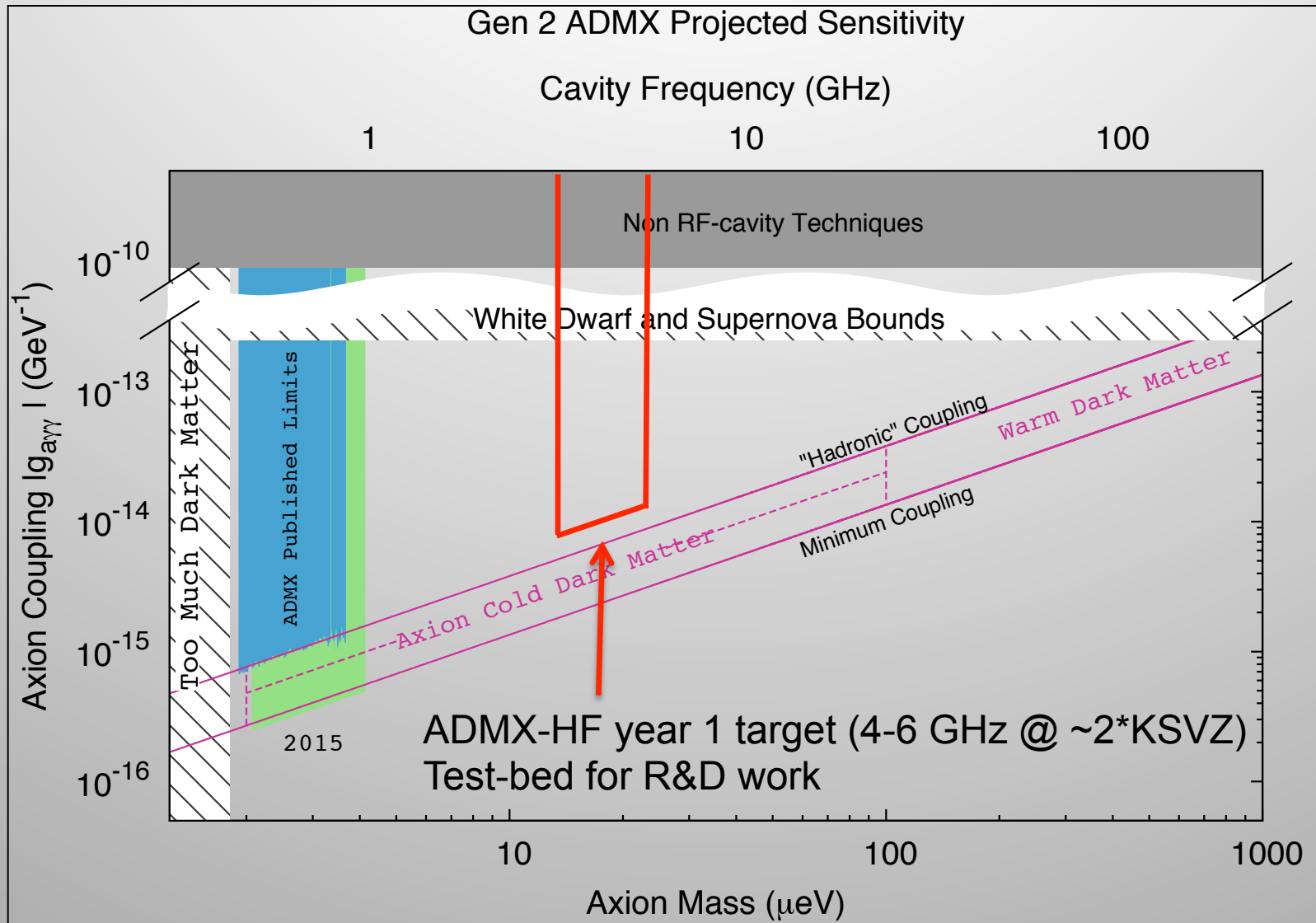


ADMX-HF Status

- Construction 2012-15
- Cavity volume $\sim 1\%$ that of ADMX
- 48 MHz run performed at $T_{\text{SYS}} \sim 2.5 T_{\text{SQL}}$
- Long run begins fall 2015 in > 5 GHz range



ADMX-HF: Year 1 prospects



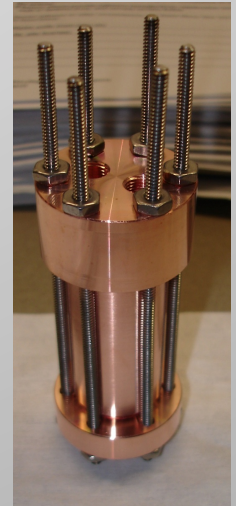
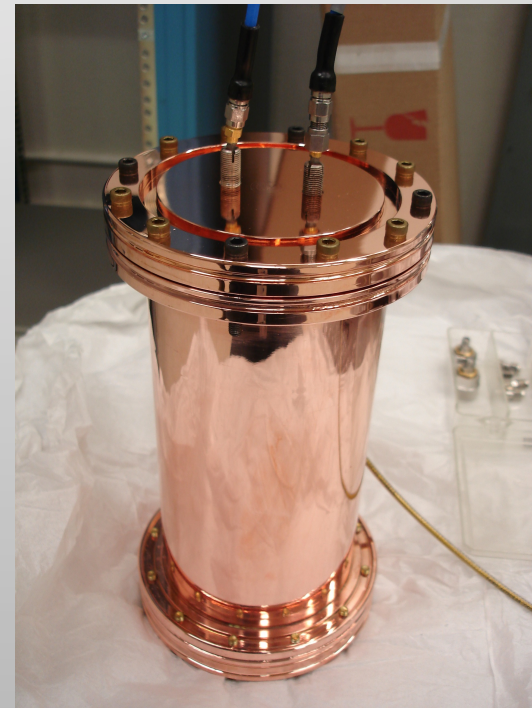
Challenges for cavities at higher frequencies/mass

- Higher Frequency requires smaller cavities – sample smaller volume!
- Quality factor goes down as frequency increases!

Radius – 19 inches
Frequency ~ 540 MHz
Q – 200,000
Axion Mass ~ 2 μeV
Volume – 220 liters

Radius – 2.5 inches
Frequency ~ 2.4 GHz
Axion Mass ~ 9 μeV
Q – 120,000
Volume ~ 2.6 liters

Radius – 0.5 inches
Frequency ~ 10 GHz
Axion Mass ~ 36 μeV
Q – 50,000
Volume – 0.025 liters



ADMX Gen2: Cavities for higher frequencies

Larger sets of tuning rods in current cavity

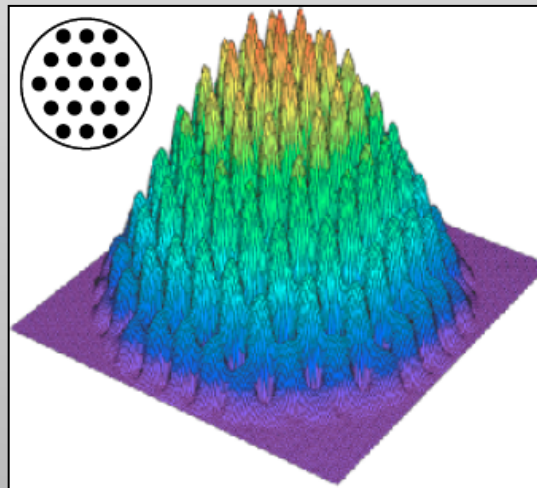
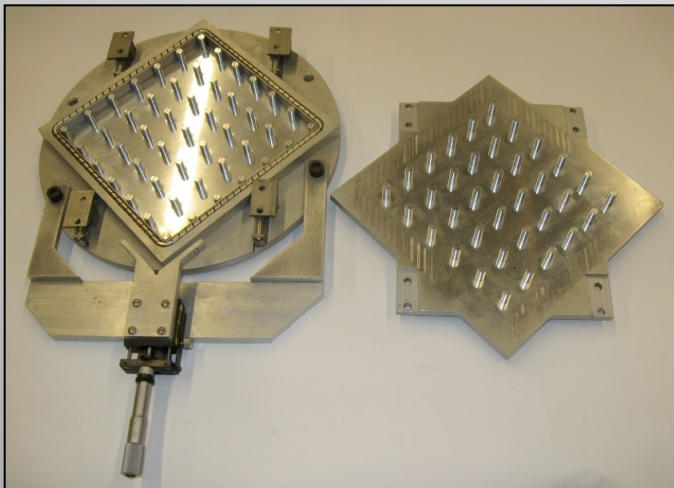
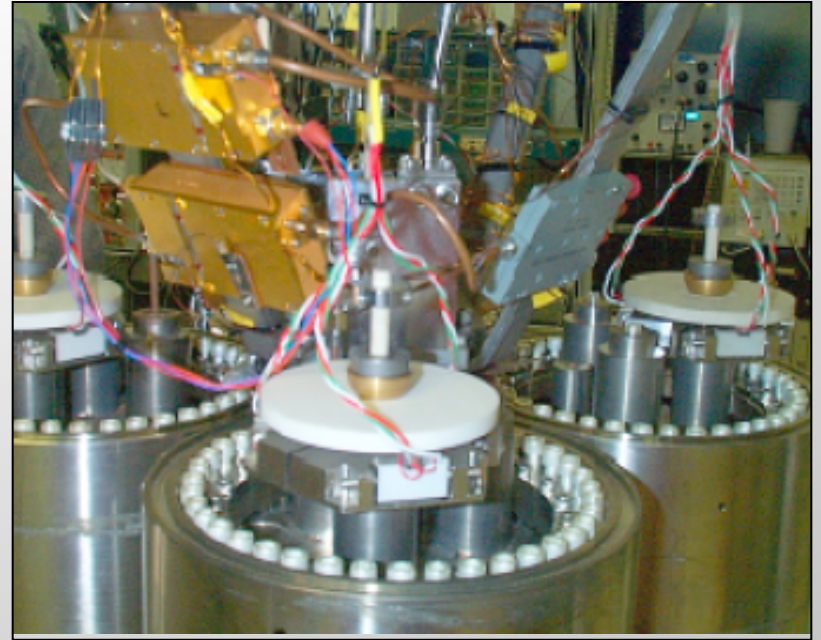
- Good to ~ 2 GHz ($Q \sim 10^5$ / $V_{ol} \sim 100$ l)

Multiple cavities in magnet bore

- Possible to use up to $\sim 3 - 4$ GHz
- Difficult to scale to > 8 cavities

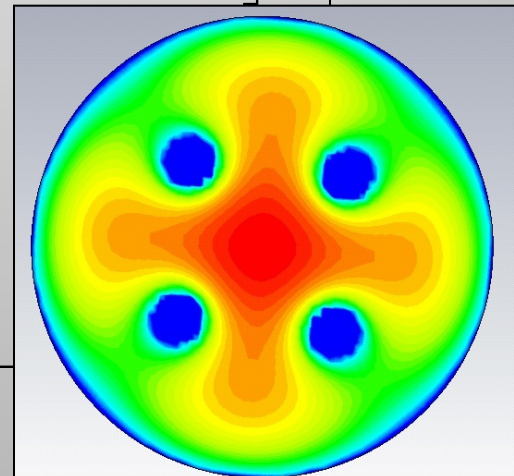
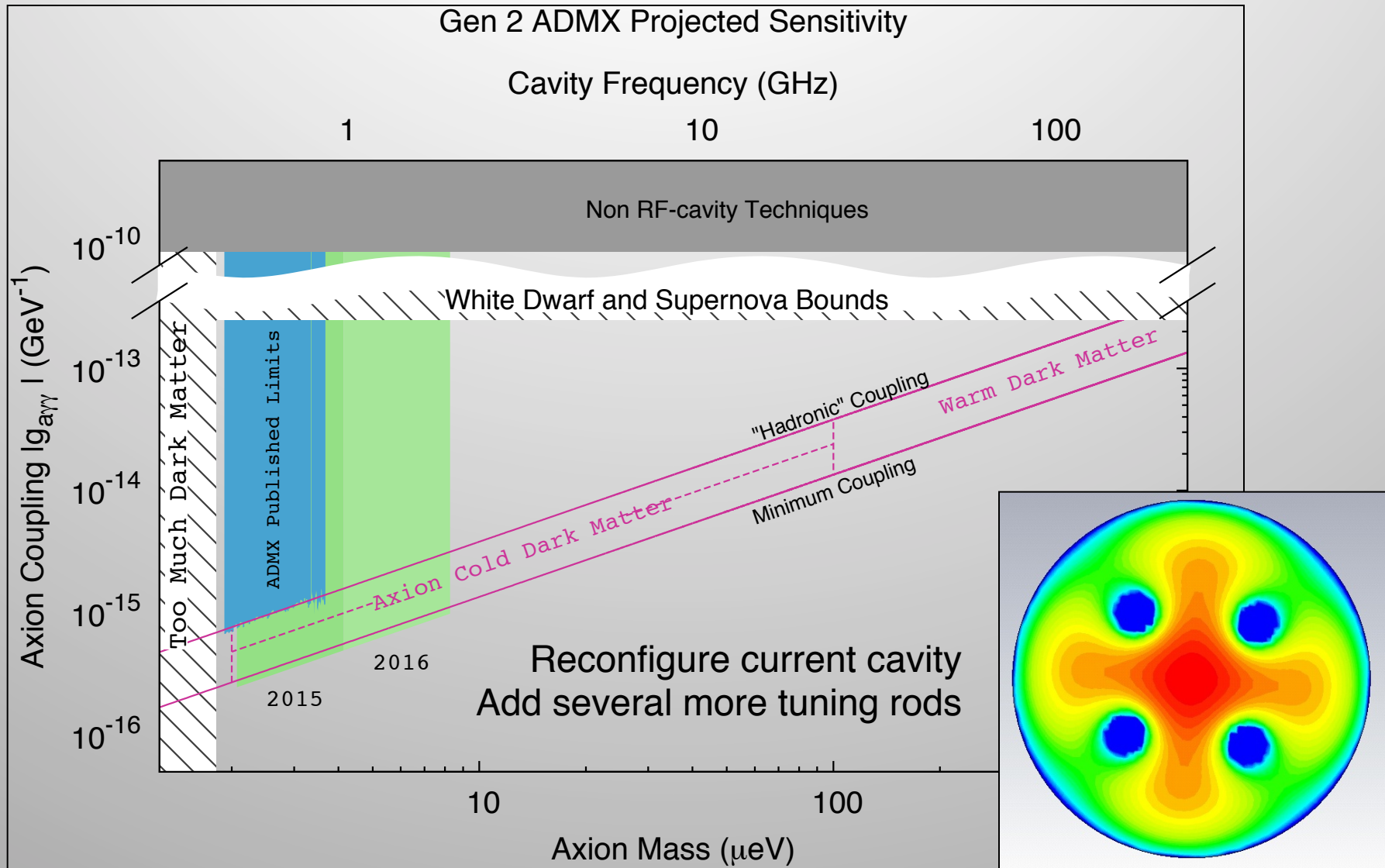
Photonic Bandgap Cavities

- Lattice of posts for higher frequencies
- Maintain relatively large E-field coverage
- Use from $4 - 10$ GHz

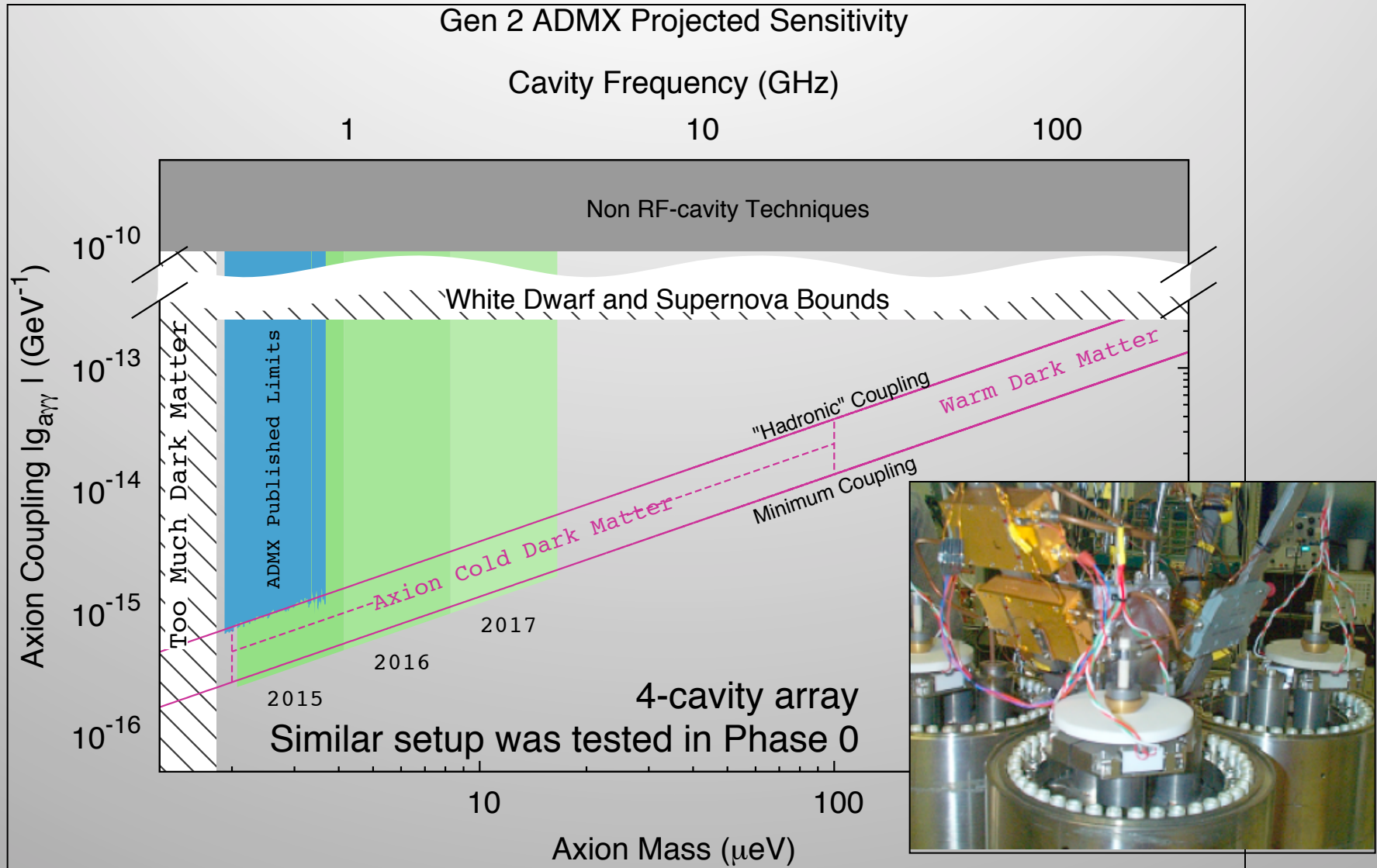


**Primary ADMX
R&D focus for the
next few years!**

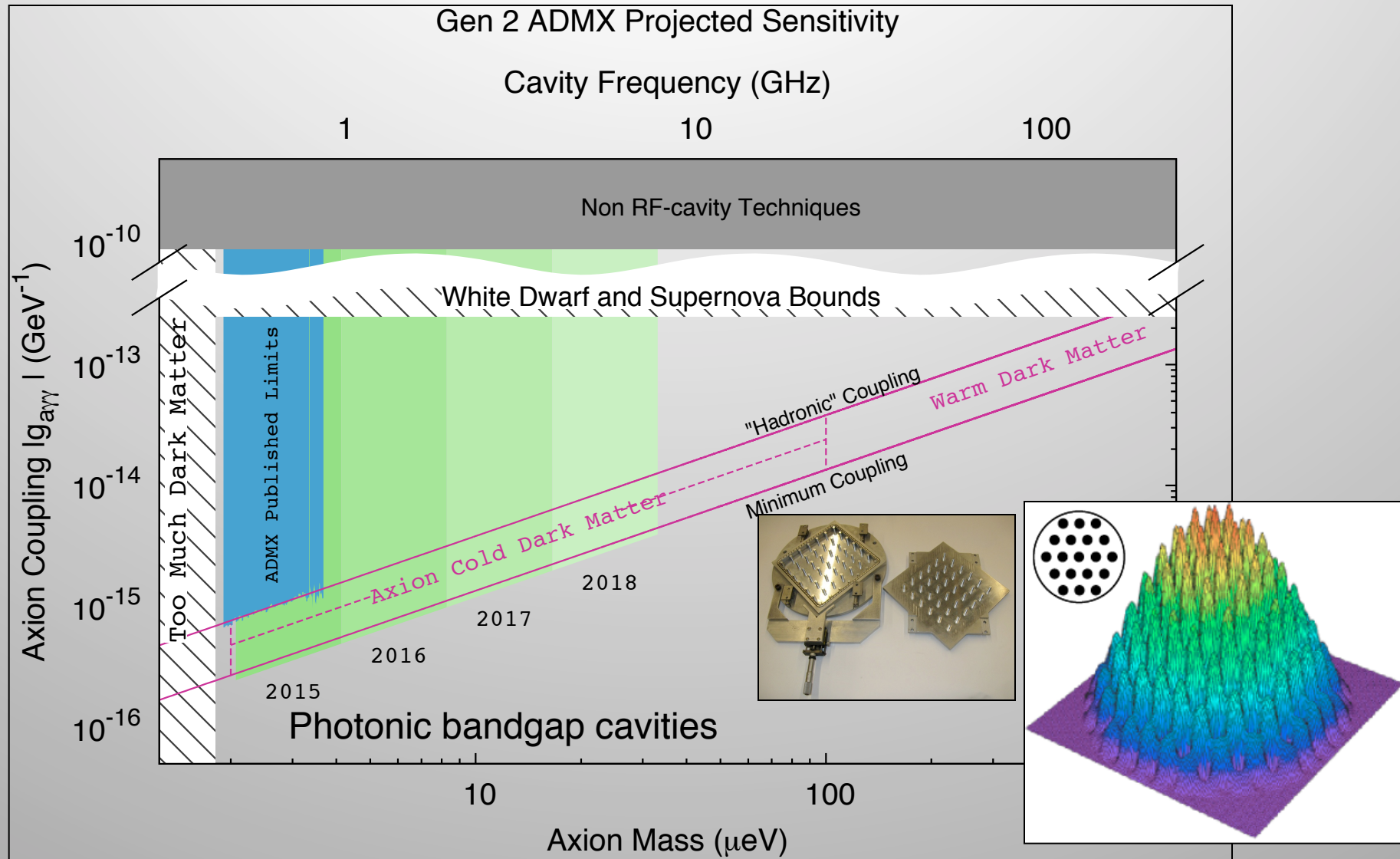
ADMX Gen 2 Science Prospects: Year 2 (1 – 2 GHz)



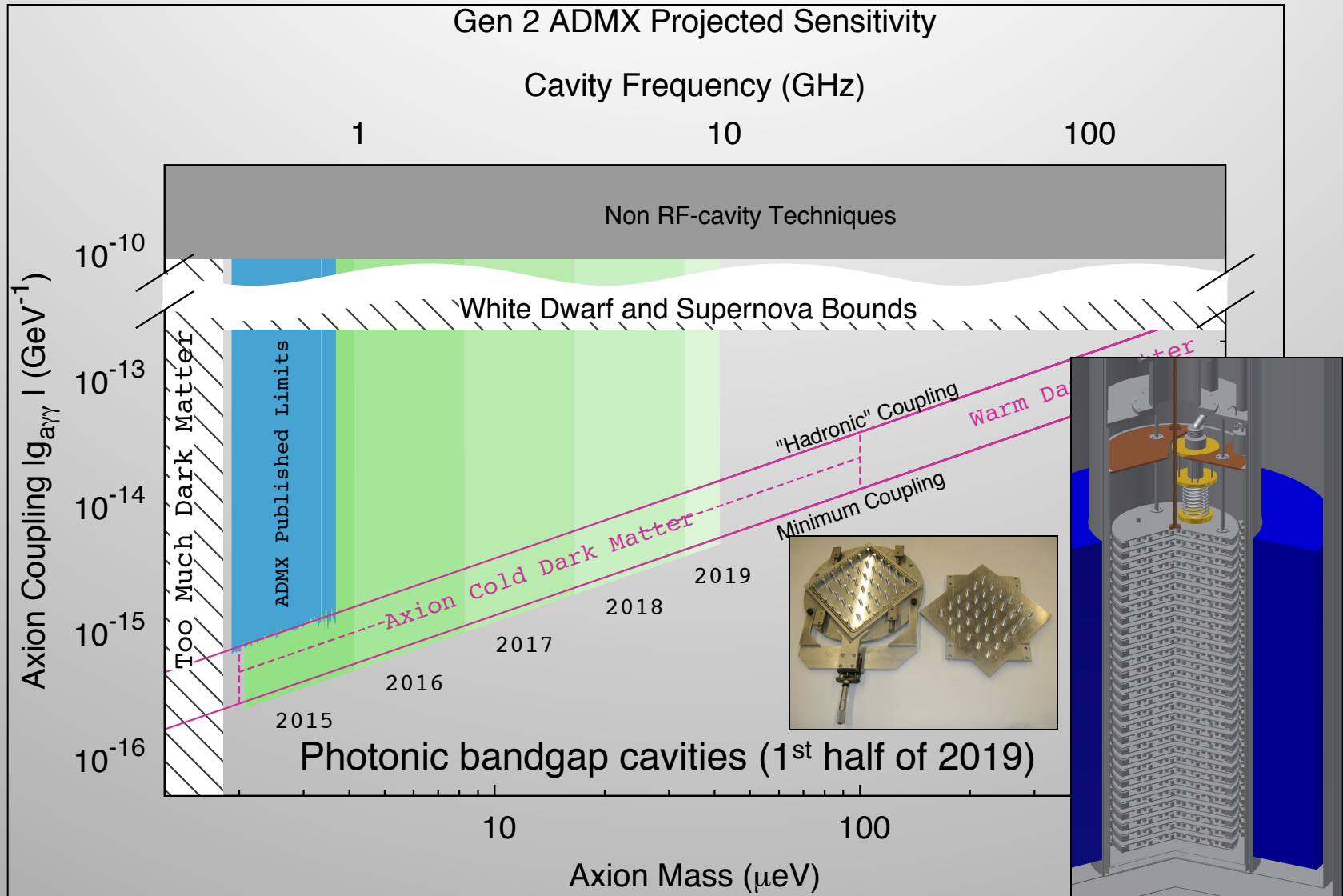
ADMX Gen 2 Science Prospects: Year 3 (2 – 4 GHz)



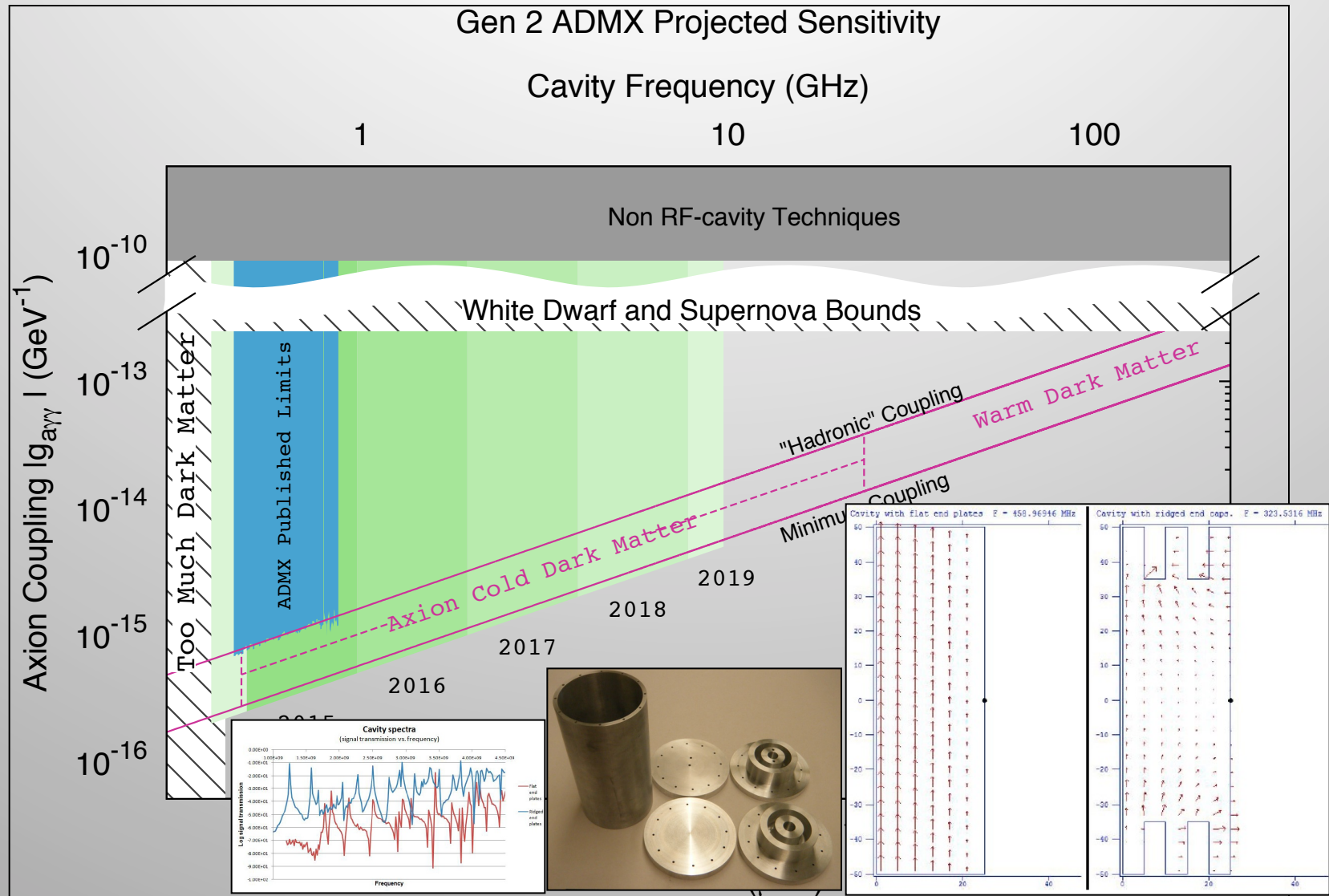
ADMX Gen 2 Science Prospects: Year 4 (4 – 8 GHz)



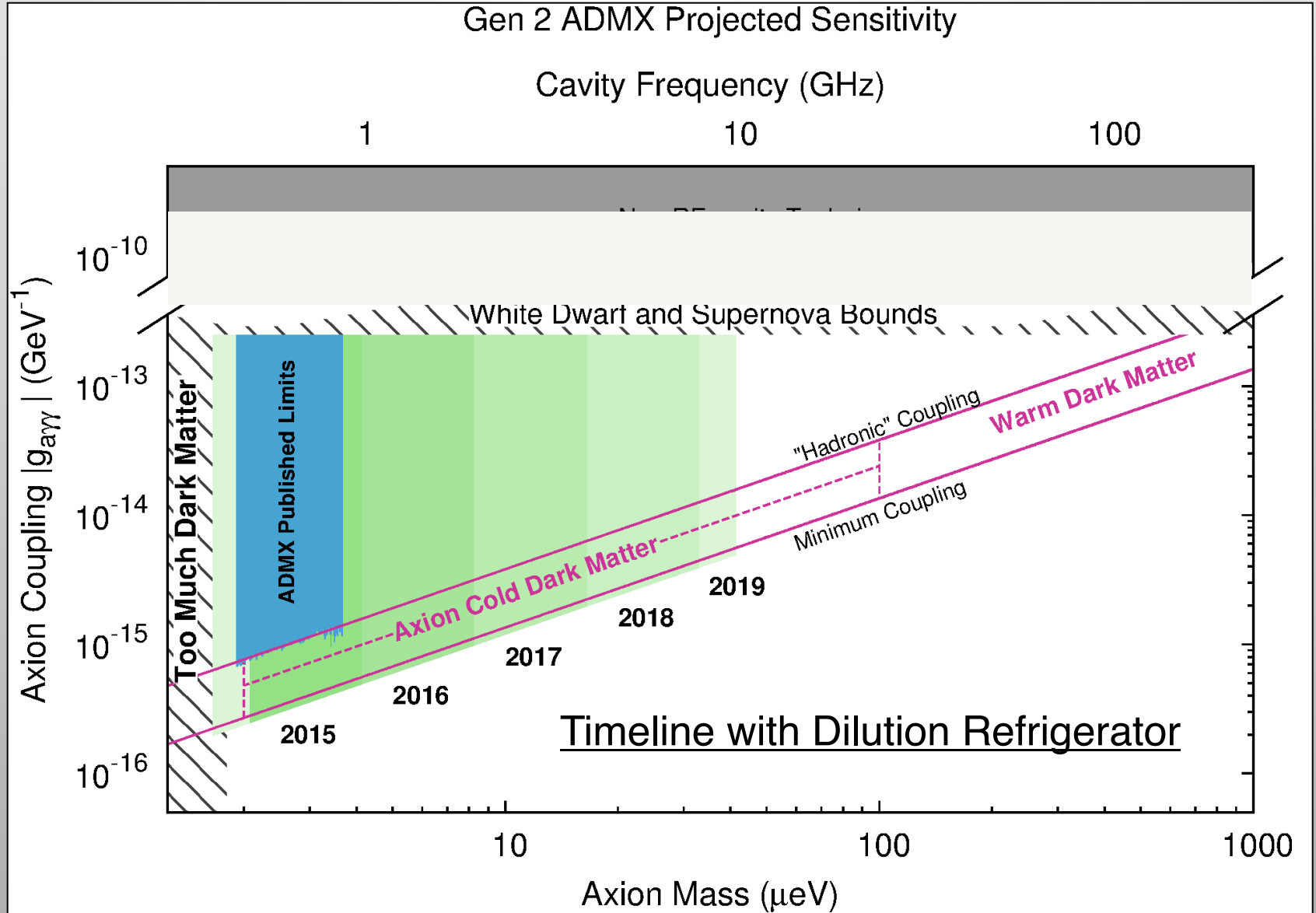
ADMX Gen 2 Science Prospects: Year 5 (8 – 10 GHz)



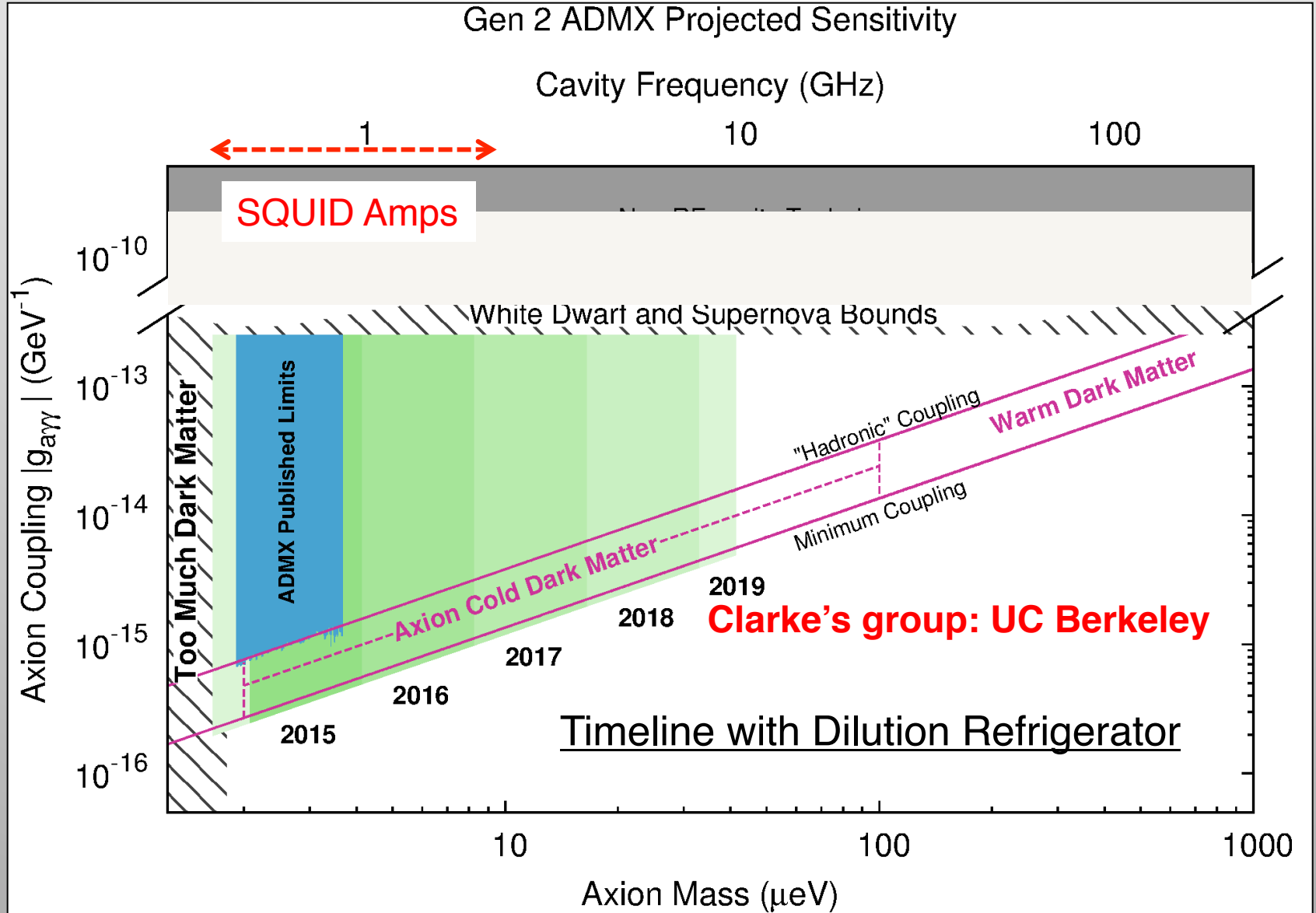
ADMX Gen 2 Science Prospects: Year 5 (< 0.5 GHz)



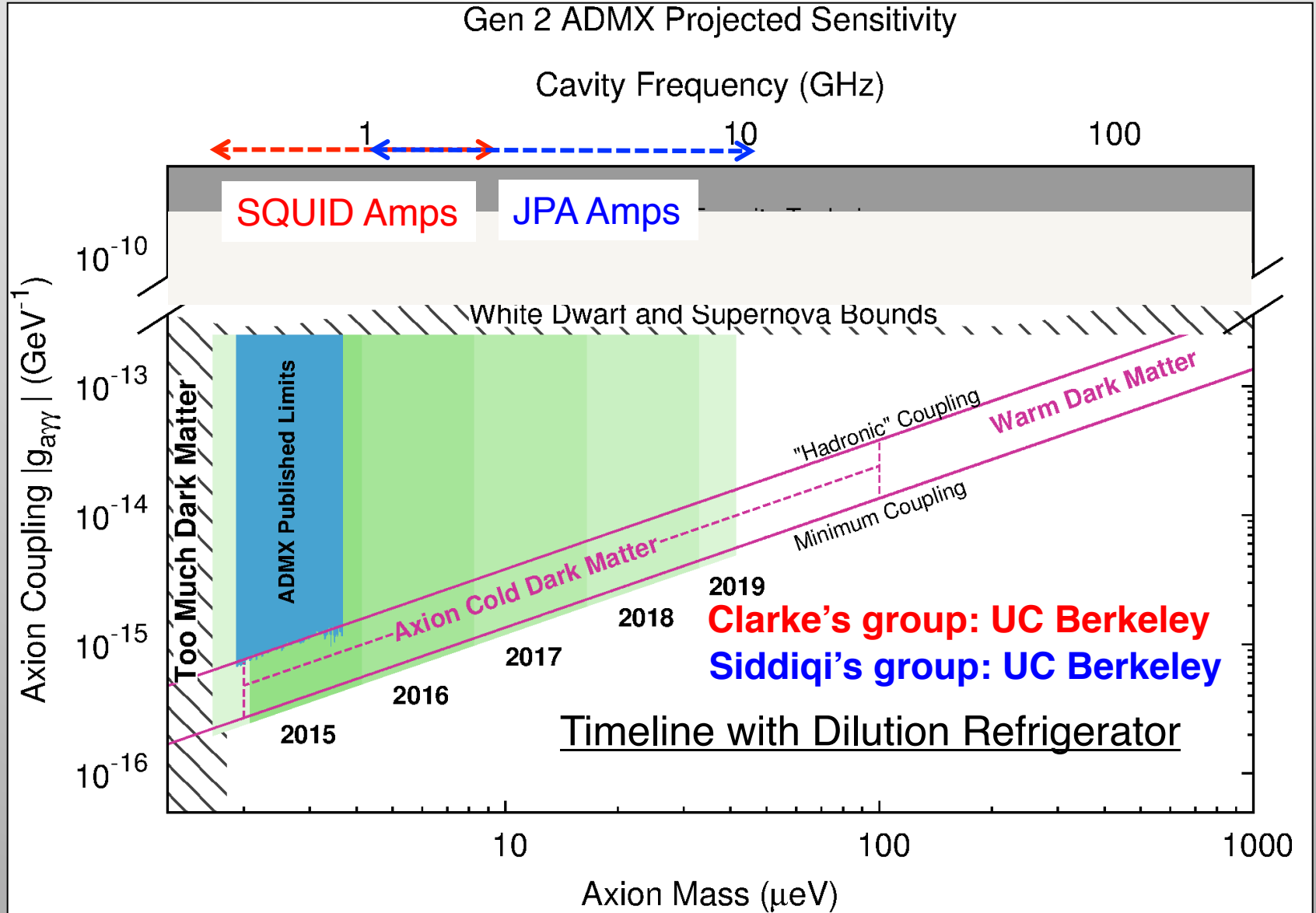
Quantum Limited Amplifiers



Quantum Limited Amplifiers



Quantum Limited Amplifiers



Going beyond the (relatively) near term searches...

New Magnet Systems: Maximize B^2V

- Magnet technology continues to improve
- Large solenoid fields (> 22 T) now possible
- Scan Rate $\sim B^4V^2$
- Prudent to invest here to broaden search
- Multiple frequencies scanned at same time

Bruker NMR magnet
23.5 T with 5.4 cm bore



ADMX magnet
8 T with 0.5 m bore



Microwave Cavity Challenges

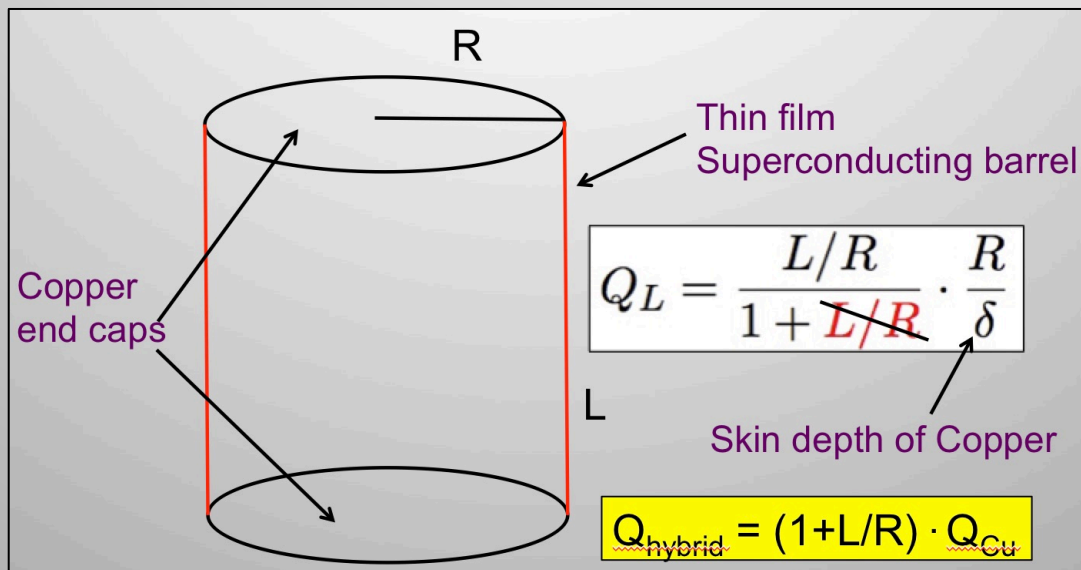
- Maximize Volume inside B-Field ($P_a \propto B^2 V$)

- Maximize Form Factor:

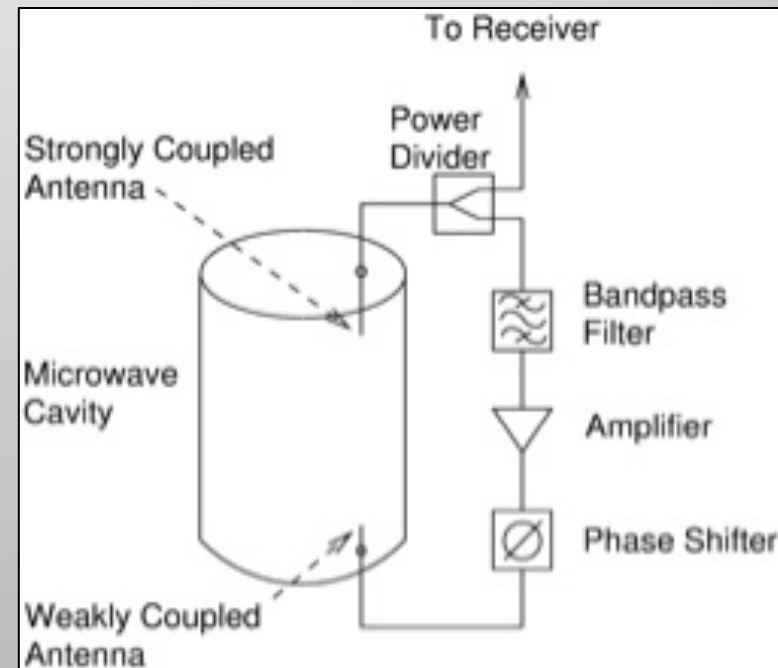
$$C_{lmn} = \frac{|\int_V d^3x \vec{E}_\omega \cdot \vec{B}_0|^2}{B_0^2 V \int_V d^3x |\vec{E}_\omega|^2}$$

- Maximize Quality Factor (up to $Q_a \sim 10^6$)

Superconducting “hybrid” cavities



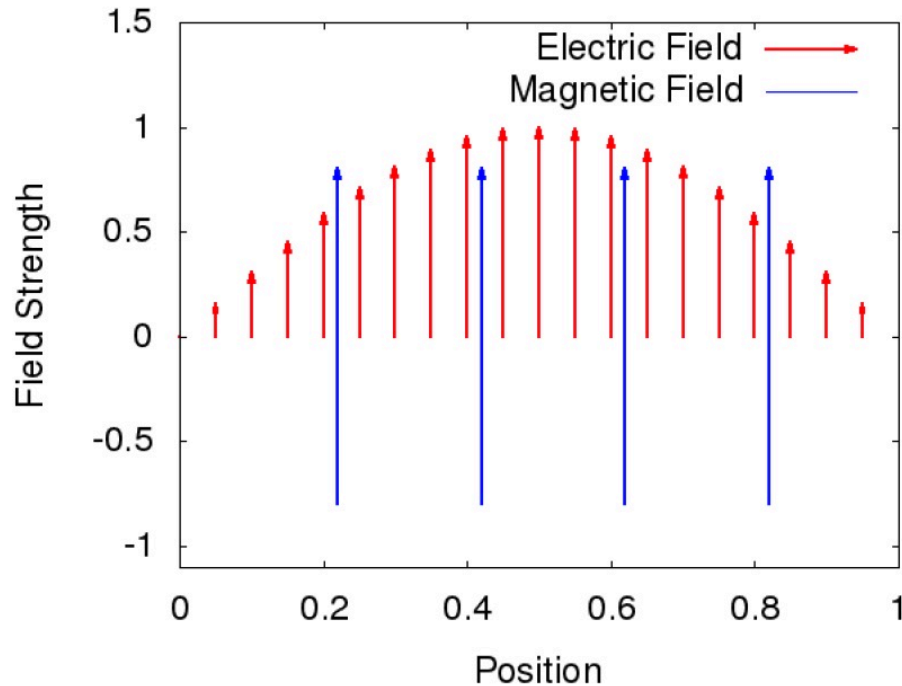
Active Feedback Cavities



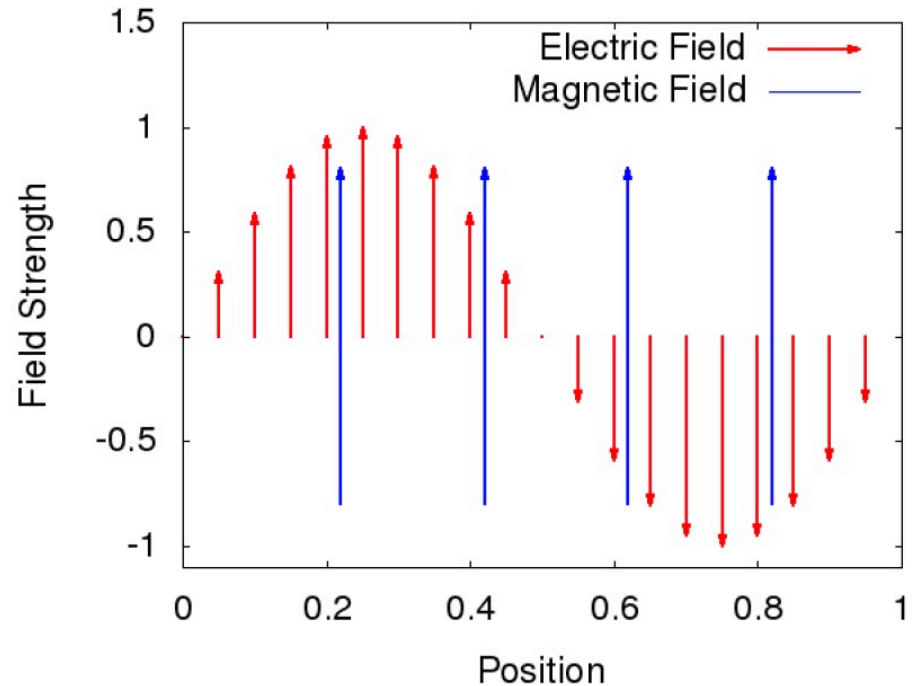
Leverage superconducting acc. cavity development

New Geometries: beyond cylinder in solenoid

Q decreases with increasing frequency, and volume decreases with wavelength in a constant magnetic field



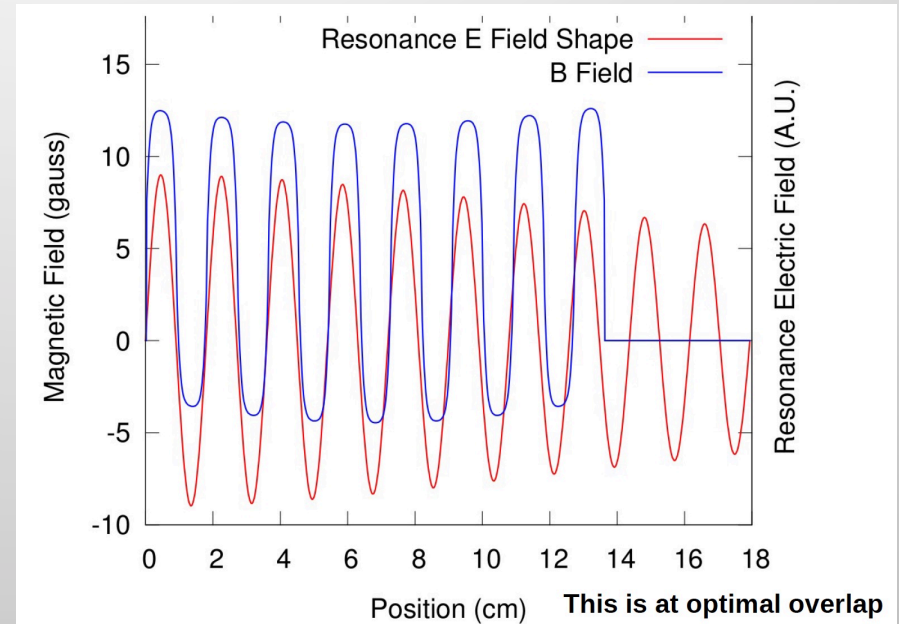
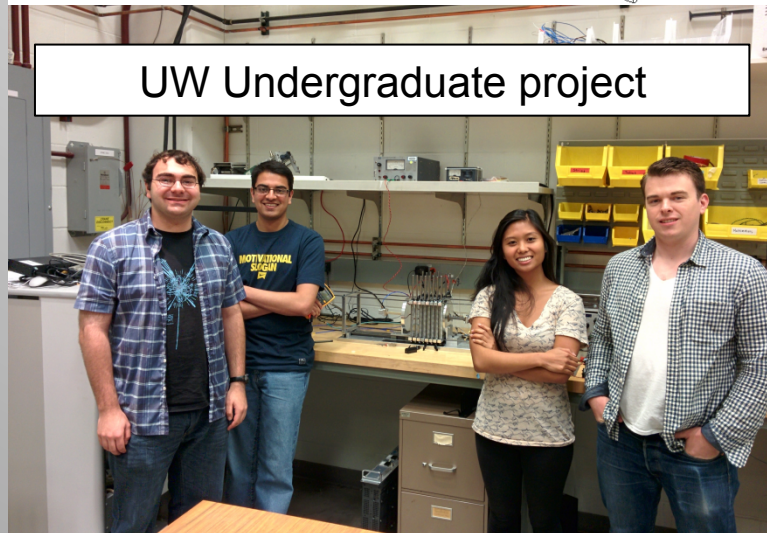
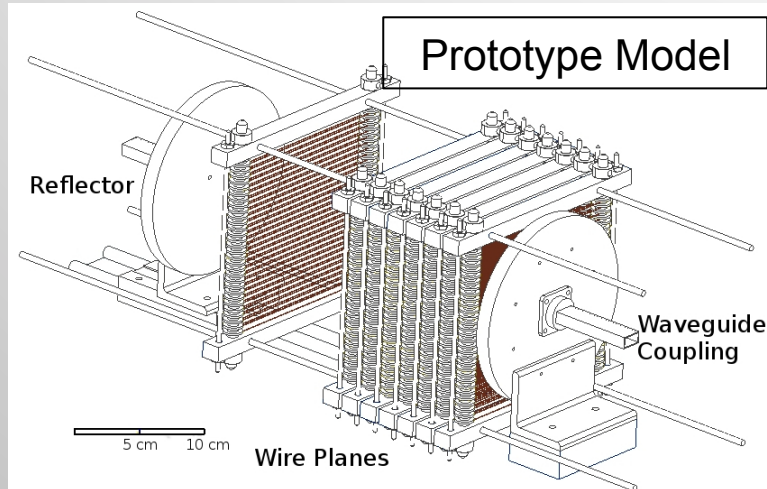
This mode couples to axions.
(Form Factor 0.8)



This mode does not.
(Form Factor 0)

New Geometries: Open Resonator R&D

Open resonators may access frequencies too high to reach with closed cavities could expand ADMX reach to highest possible dark matter axion masses

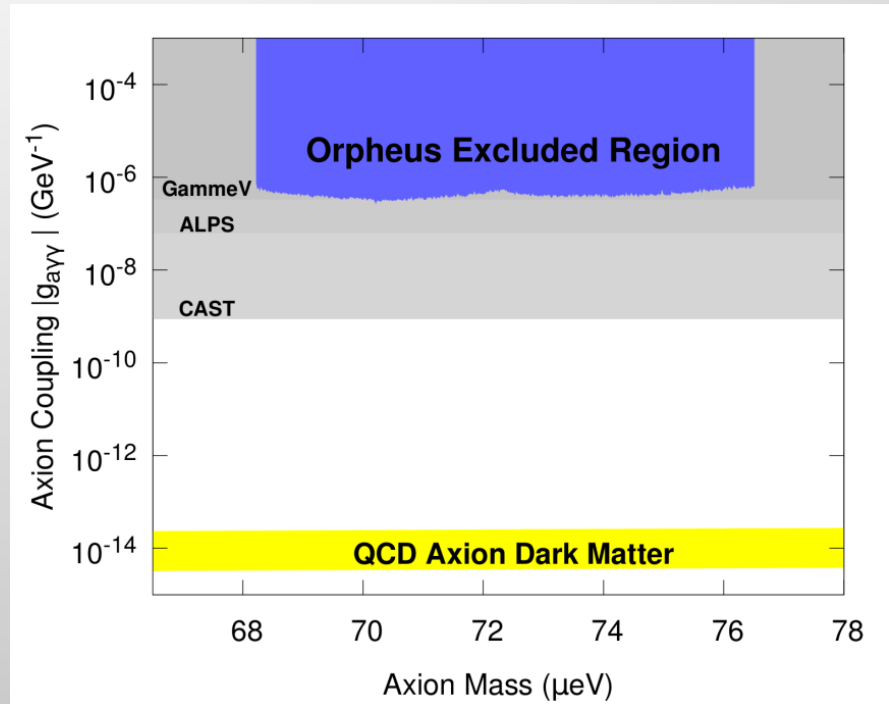
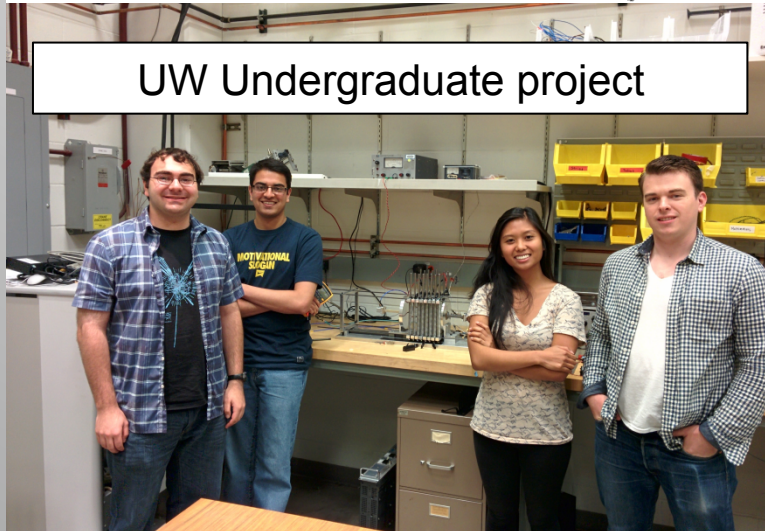
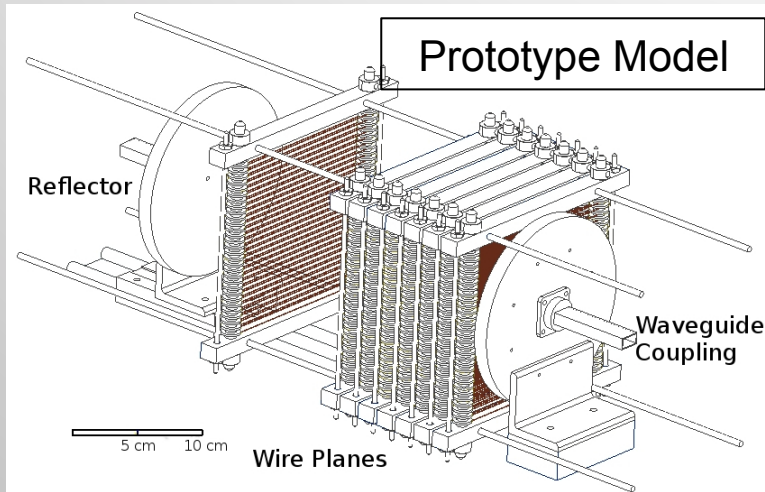


PhysRevD.91.011701

System potentially good to much higher frequencies (40 GHz or more)

Open Resonator R&D

Open resonators may access frequencies too high to reach with closed cavities could expand ADMX reach to highest possible dark matter axion masses

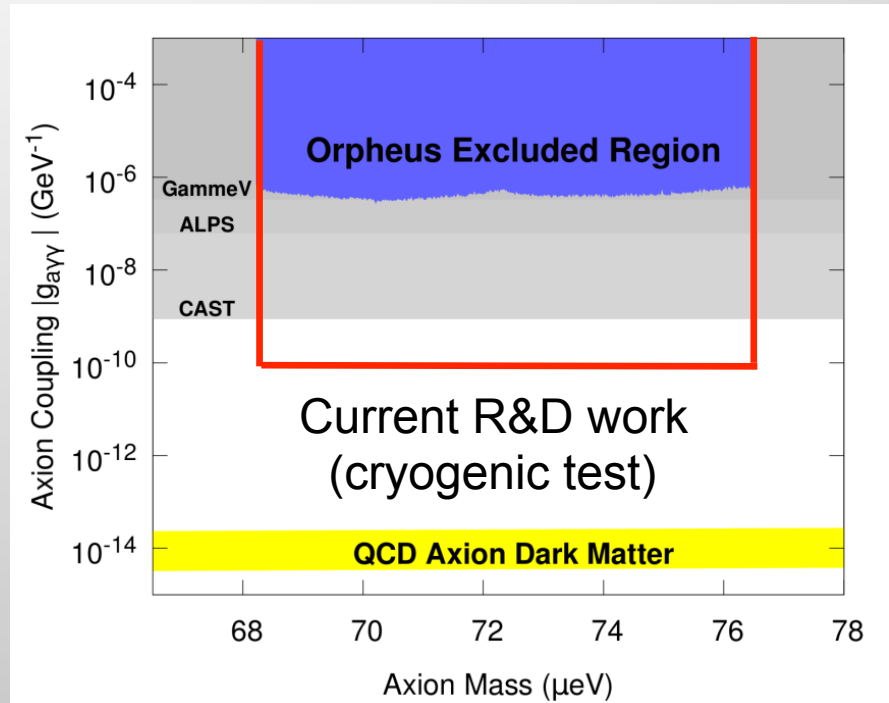
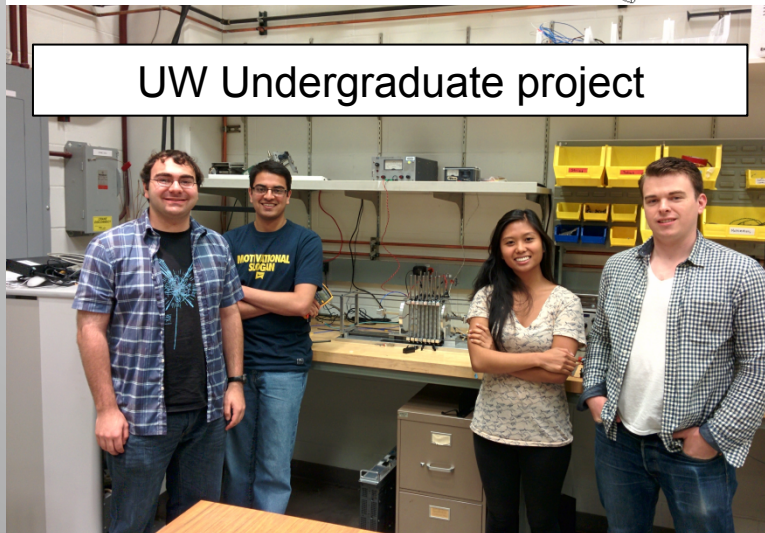
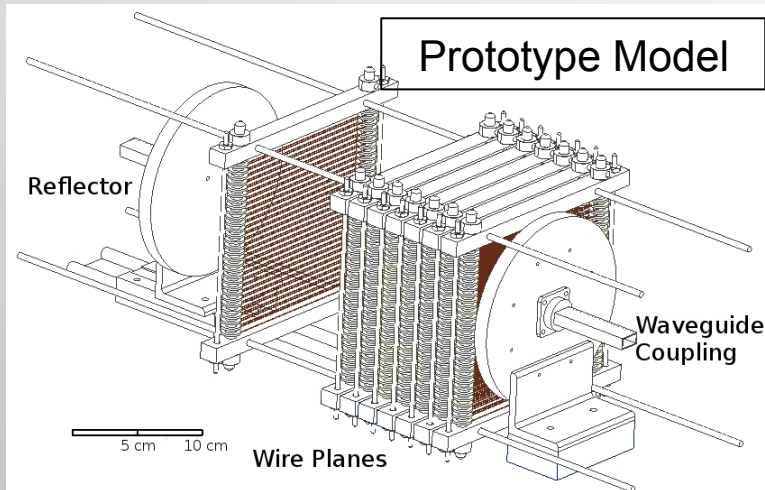


PhysRevD.91.011701

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Open Resonator R&D

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PhysRevD.91.011701

System potentially good to much higher frequencies (40 GHz or more)

Linear amplifiers are subject to the Standard Quantum Limit

$$T_N > T_{SQL} \quad \text{where} \quad k_B T_{SQL} = h\nu$$

ν [GHz]	m_a [μeV]	T_{SQL} [mK]
0.5	2.1	24
5	20.7	240
20	82.8	960

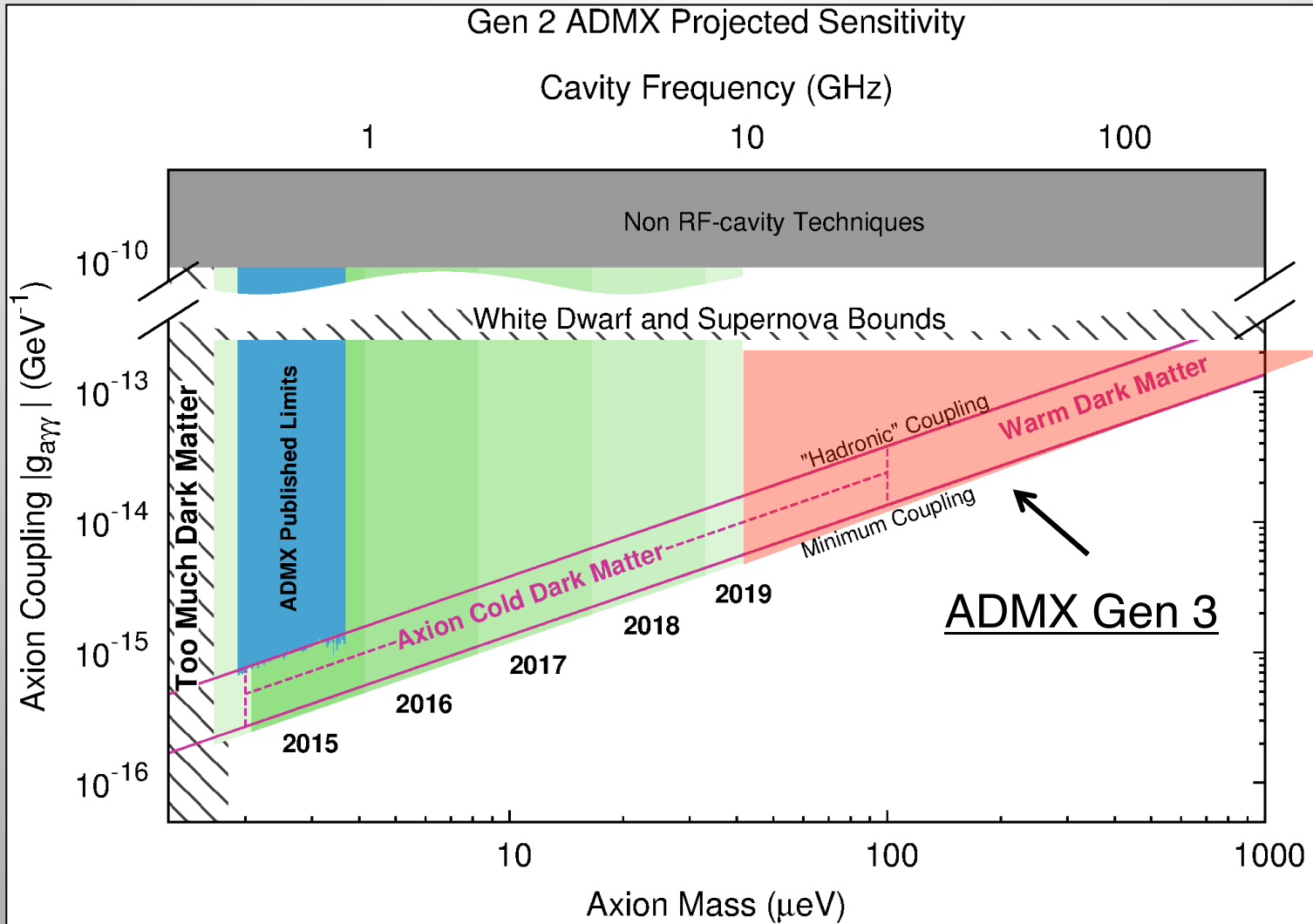
The SQL can be evaded by:

- Squeezed-vacuum state receiver (e.g. GEO, LIGO)
- Single-photon detectors (e.g. qubits, bolometers)

* S.K. Lamoreaux et al. (PhysRevD.88.035020)

ADMX Generation 3

Long term goal is to detect or rule out axion as primary dark matter candidate.

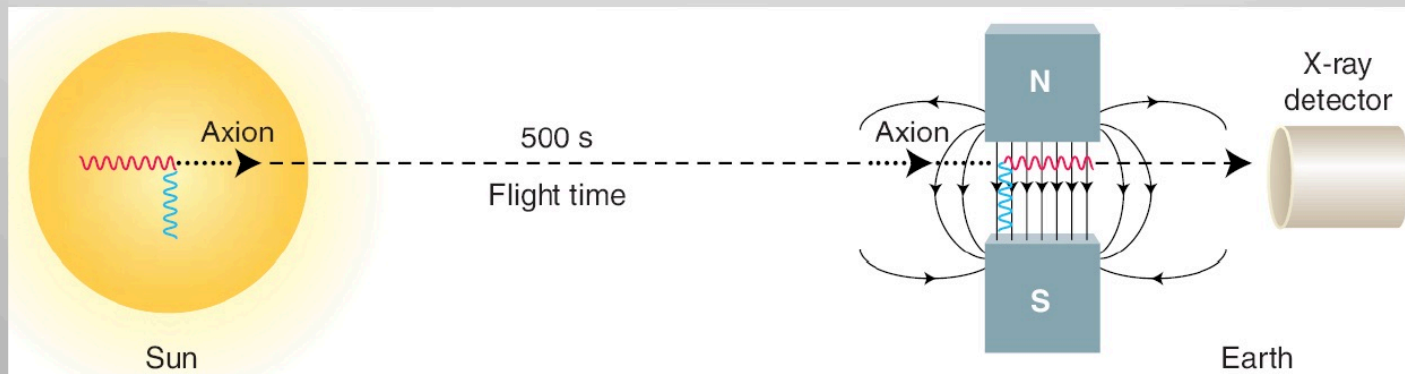


Production and detection of axions

- First axion helioscope proposed by P. Sikivie

Sikivie *PRL* 51:1415 (1983)

- Blackbody photons (keV) in solar core can be converted into axions in the presence of strong electromagnetic fields in the plasma
- Reconversion of axions into x-ray photons is possible in strong laboratory magnetic fields

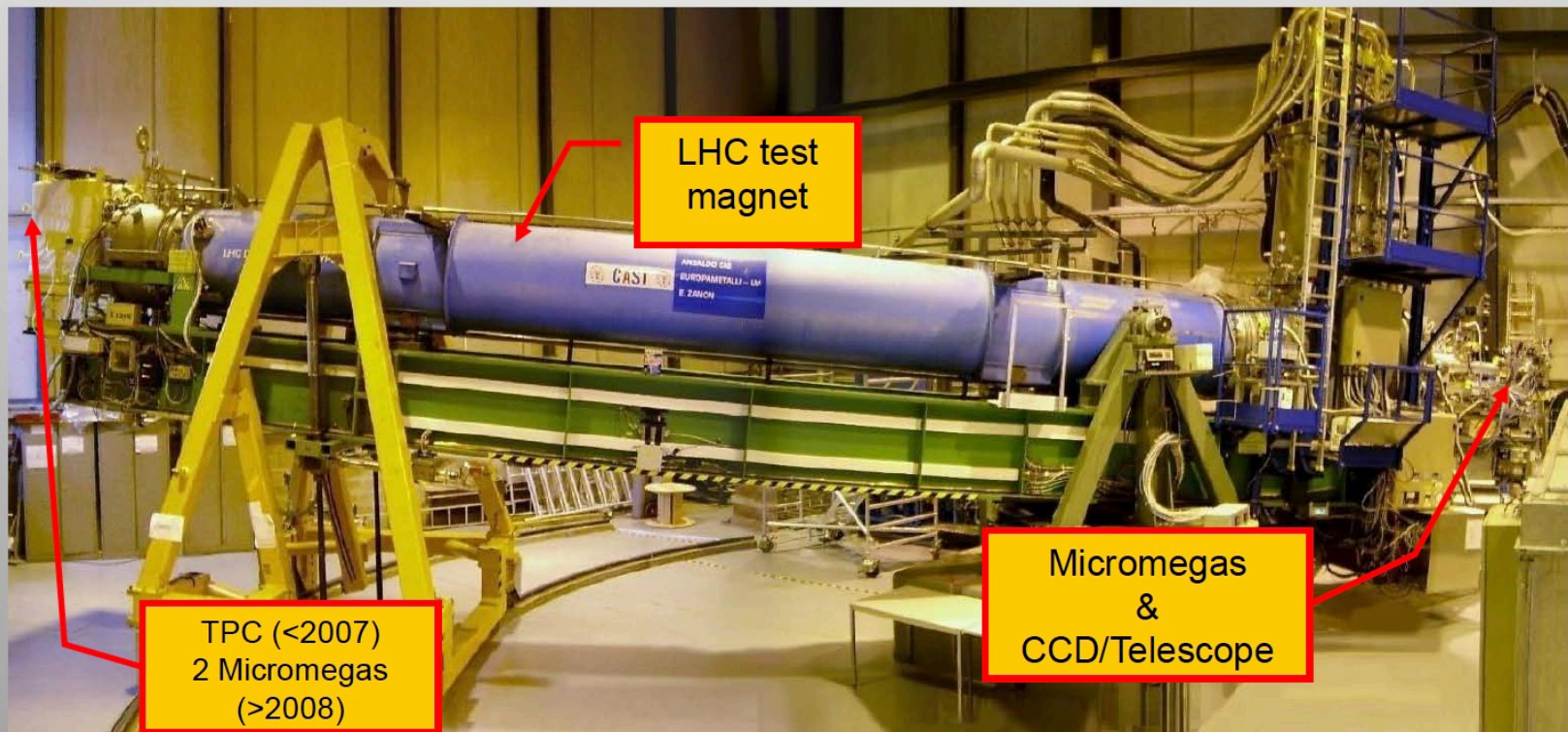


- Idea refined by K. van Bibber et al. by using buffer gas to restore coherence over long magnetic field

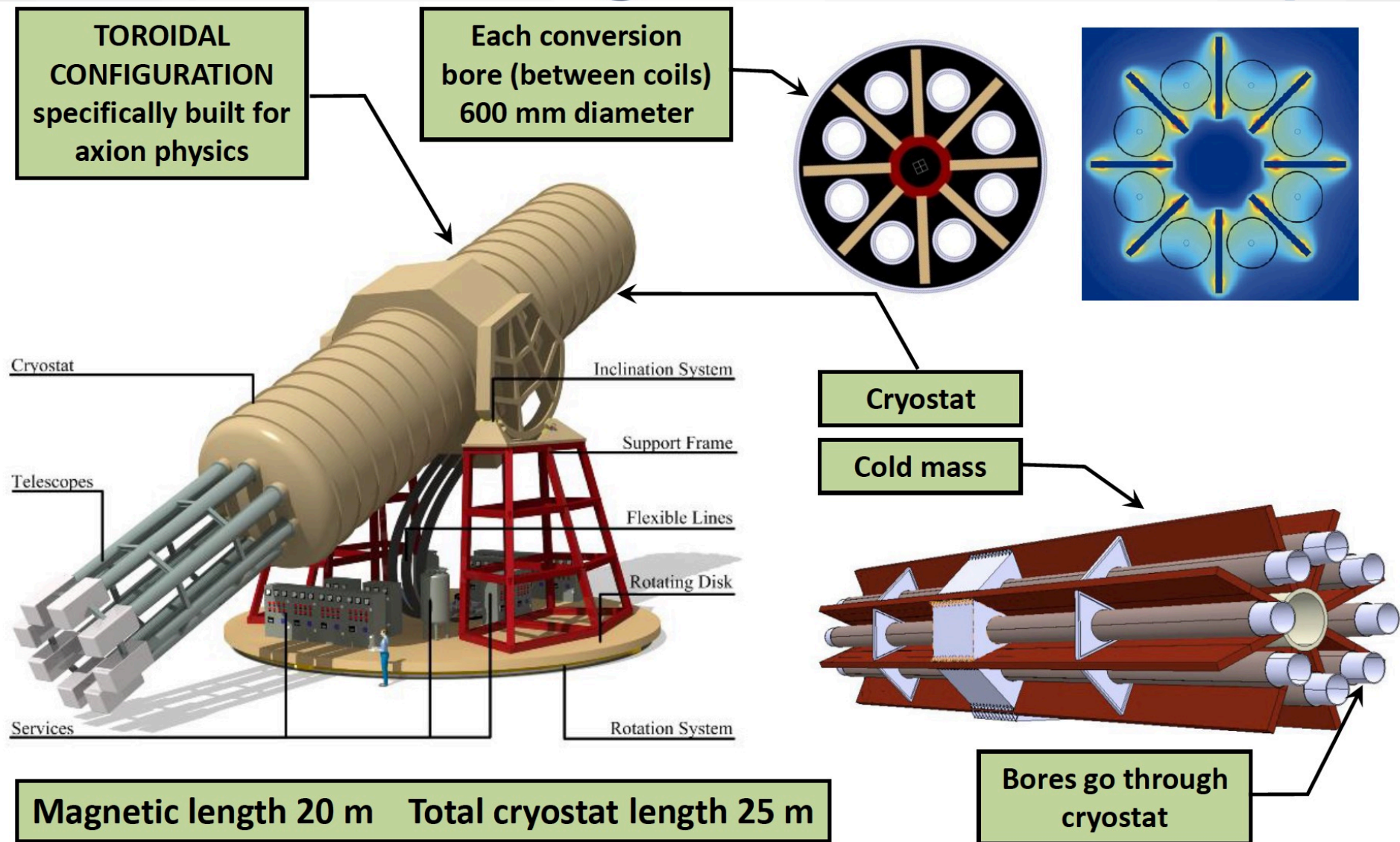
Van Bibber et al. *PhysRevD* 39:2089 (1989)

CAST experiment @ CERN

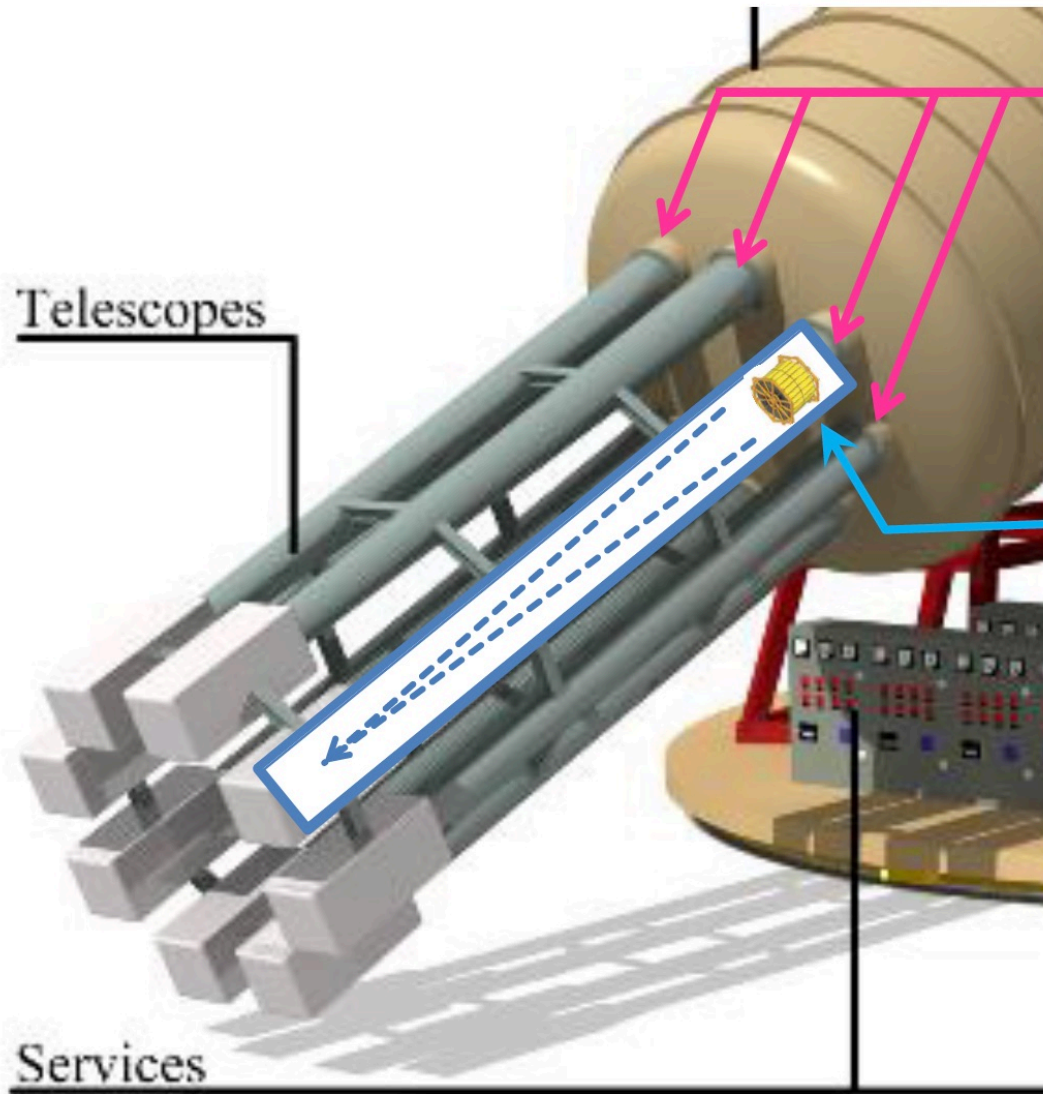
- International collaboration started in 1999
- Almost continuously acquired data since 2003
- 20 institutions from 11 countries, approximately 70 PhD scientists
- Thesis project for 10 PhD students, 6 more pending
- **Very mature technology → CAST is 3rd generation helioscope**



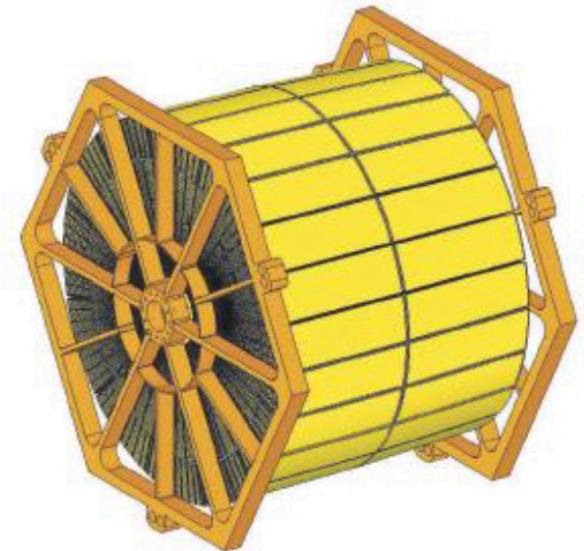
IAXO – The new generation helioscope



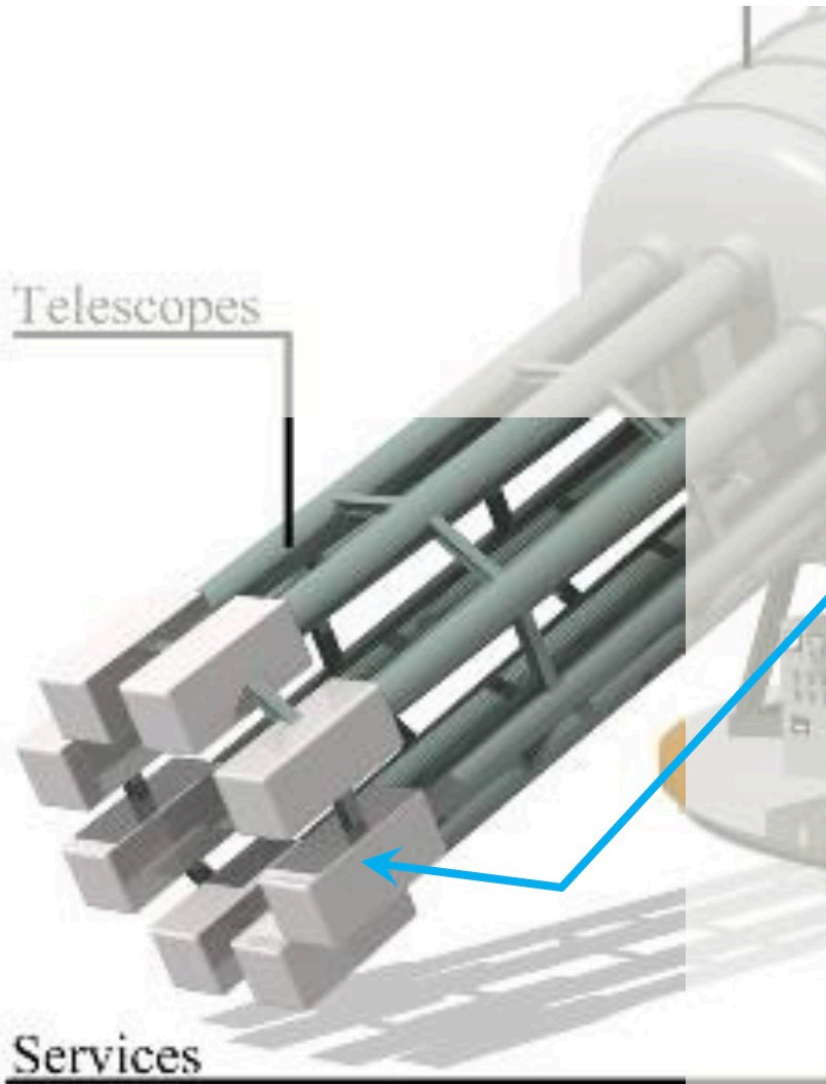
Solar axions energy primarily kinetic (keV X-rays)



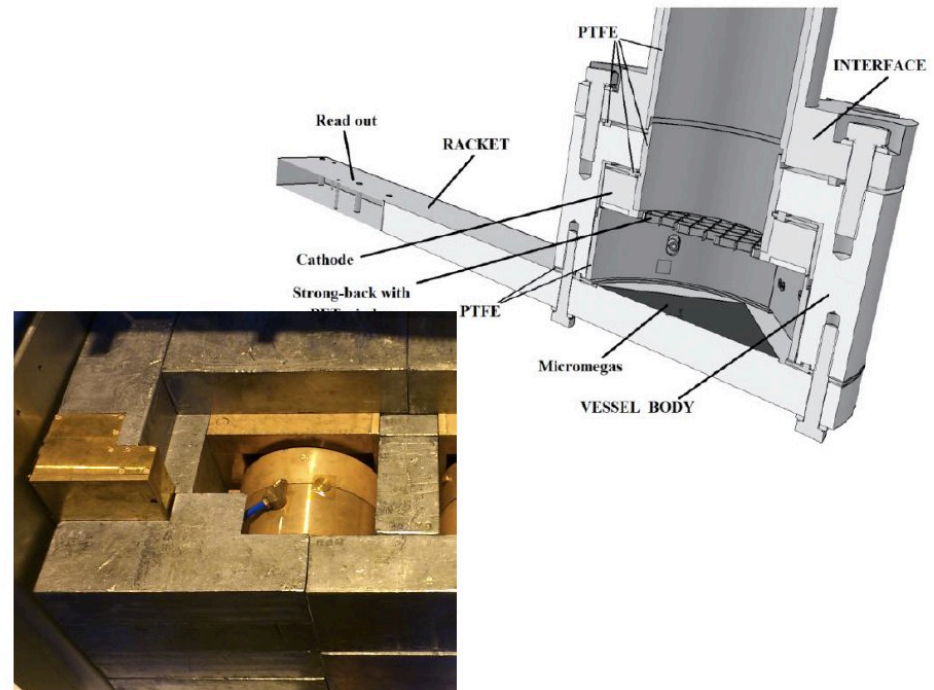
- Each bore has an x-ray telescope
- Exquisite imaging not required
- Need cost-effective way to build 8 (+1 spare), highly nested optics
- Baseline: Use approach from NASA's *NuSTAR* satellite



Solar axions energy primarily kinetic (keV X-rays)



- 8 detector systems
- Small gas chamber with Micromegas readouts for low-background x-ray detection
- Optimized shielding



Sensitivity of CAST vs IAXO

**MAGNET &
TRACKING**

**X-RAY
OPTICS**

**X-RAY
DETECTORS**

$$g_{a\gamma}^4 \propto$$

$$\underbrace{(BL)^{-2} A^{-1}}_{\text{magnet}} \times \underbrace{t^{-1/2}}_{\text{exposure}}$$

B = magnetic field

L = magnet length

A = cross-sectional area

t = time

\times

$$\underbrace{s^{1/2} \varepsilon_0^{-1}}_{\text{optics}}$$

s = spot size

ε_0 = efficiency

\times

$$\underbrace{b^{1/2} \varepsilon^{-1}}_{\text{detectors}}$$

b = background

ε = efficiency

$\geq 1050 \times$ better

$1.4 \times$ better

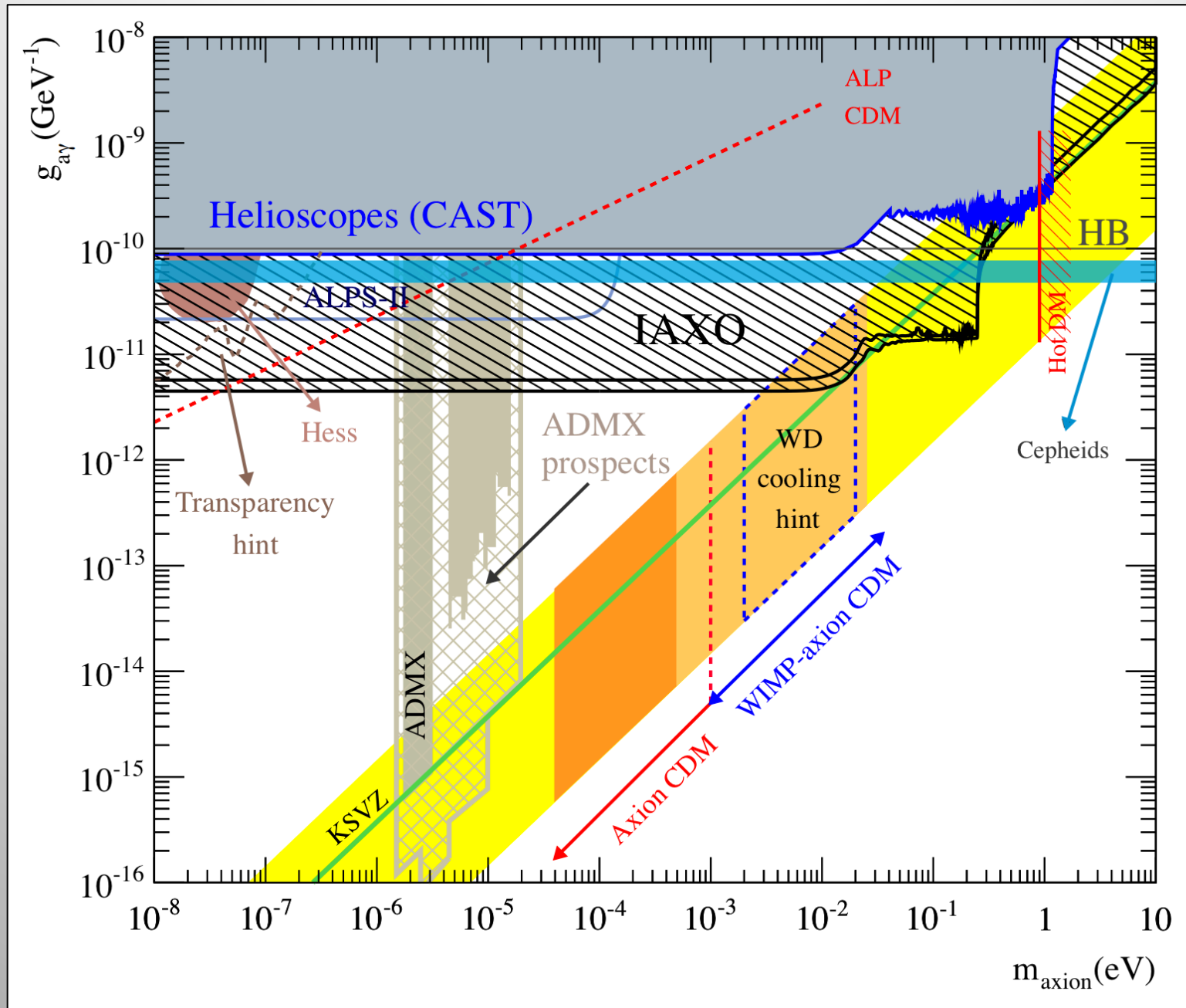
$14 \times$ better

IAXO, relative to CAST:

$$S/N = g_{a\gamma}^4 \geq 20000$$



$$g_{a\gamma} \geq 12 \times \text{better}$$



Summary & Conclusions

Axions: solve the Strong-CP problem and are a compelling DM candidate

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- A narrow band experiment with concurrent R&D
Takes data in one mass range while developing systems for higher masses.

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Technologies that are under active development (haloscopes):

1. Microwave Cavities:
 - High-Frequency, Large-Volume Tunable Systems with high Q
2. RF Detectors: Quantum Limited (0.25 – 10 GHz): SQUIDs & JPAs
3. Beyond several GHz the standard quantum limit begins to dominate
 - Employ Squeezed States and Eventually Single-Photon-Counters
4. Large Magnets can increase axion conversion signal.

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Stay tuned for additional talks by Jeremy Mardon & Surjeet Rajendran